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**A RADIATIVE HEATING AND COOLING
ALGORITHM FOR A NUMERICAL MODEL
OF THE LARGE SCALE STRATOSPHERIC
CIRCULATION**

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1. INTRODUCTION

A radiative heating and cooling algorithm has been developed for use in the dynamical model of the middle atmosphere written by J. R. Holton (Holton and Wehrbein, 1979; Holton and Wehrbein, 1980). A Curtis matrix is used to compute cooling by the 15μ and 10μ bands of carbon dioxide. Escape of radiation to space and exchange with the lower boundary are used for the 9.6μ band of ozone. Voigt line shape, vibrational relaxation, line overlap, and the temperature dependence of line strength distributions and transmission functions are incorporated into the Curtis matrices. Properties of this algorithm have been outlined in Wehrbein and Leovy (1981). This report is a more detailed description of its development.

Section 2 discusses the distributions of the atmospheric constituents included in this algorithm. Section 3 describes the method used to compute the Curtis matrices. Cooling or heating by the 9.6μ band of ozone is discussed in section 4. A description of the FORTRAN programs and sub-routines that have been developed in this study is given in section 5. Listings of these routines are found in the appendix.

2. ATMOSPHERIC CONSTITUENTS

With the exception of molecular oxygen, which absorbs solar radiation in the upper mesosphere, the atmospheric constituents that determine the thermal structure of the middle atmosphere are exceedingly minor constituents. A discussion of the contributions of each gas is found in London (1980). In the present work we shall consider only carbon dioxide (15μ and 10μ bands), and ozone (9.6μ band and solar absorption), and molecular oxygen (solar absorption). A more sophisticated algorithm might include the overlap of carbon dioxide and ozone bands with each other, and with the bands of tropospheric water vapor when computing radiative exchange with the ground (e.g., Ramanathan, 1976).

2.1 Carbon dioxide

Carbon dioxide is, for our purposes, mixed throughout the atmosphere with volume mixing ratio of 0.33×10^{-3} . Therefore the amount of CO_2 between any two pressure levels is always the same, and the transmission between two pressure levels depends only on the temperature of the intervening material.

2.2 Ozone

Both the local density and the column abundances of ozone and molecular oxygen are required to compute the radiative heating due to absorption of solar radiation. The specific ozone density $Q_3 \equiv n_3/n_{\text{STP}}$ where $n_{\text{STP}} = 2.6869 \times 10^{19} \text{ cm}^{-3}$ and n_3 is the number density of ozone. The mass mixing ratio (in grams/gram) is given by $W_3 = \rho_3/\rho_a$, where ρ_3 and ρ_a are the mass densities of ozone and air, respectively. The specific density is related to the mass mixing ratio by

$$Q_3 = \frac{p}{R_a T} \frac{W_3}{m_3 n_{STP}} \quad (2.1)$$

where p and T are local pressure and temperature, R_a is the gas constant for air, and m_3 is the mass of a single ozone molecule.

The specific column abundance (in cm at STP) is given as

$$\eta_3 = \int_z^\infty Q_3(z') dz' = \frac{1}{m_3 n_{STP} g} \int_0^p W_3(p') dp' \quad (2.2)$$

Above a certain altitude z_1 (pressure level p_1) we assume that ozone is distributed with a constant scale height H_3 so that

$$\eta_3(z_1) = \frac{1}{n_{STP}} \int_z^\infty n(z_1) e^{-(z-z_1)/H_3} dz = Q_3(z_1) H_3$$

Below p_1 the mean value of W_3 in each layer is used to obtain

$$\eta_3(p_J) = \eta_3(p_1) + \frac{1}{m_3 n_{STP} g} \left\{ \sum_{j=1}^{J-1} \frac{1}{2} [W_3(p_j) + W_3(p_{j+1})] (p_{j+1} - p_j) \right\}$$

For arbitrary p η_3 is interpolated linearly between levels:

$$\eta_3(p) = \eta_3(p_J) \gamma + (1 - \gamma) \eta_3(p_{J+1}) \quad \text{where} \quad \gamma = \frac{p - p_J}{p_{J+1} - p_J} \quad (2.3)$$

for $p_J \leq p \leq p_{J+1}$.

2.3 Molecular oxygen

Assume that the mixing ratio of oxygen in air is 0.2095 up to 155 km, (3.645 x 10⁻⁹ atmospheres), above which oxygen follows the profile

$$n_2(z) = 1 \times 10^{17} e^{-z/30\text{km}} .$$

Below 155 km $n_2(z) = 0.2095 p/k_B T$ where k_B is Boltzmann's constant.

The column abundance below 155 km is therefore given by

$$\eta_2 = \int_0^z n_2(z') dz' = 1.711 \times 10^{15} + \frac{0.2095}{m_a g} (p - 3.645 \times 10^{-9} p_0) \quad (2.4)$$

where m_a is the mean mass of a molecule of air, p_0 is standard surface pressure, and η_2 is given in molecules cm⁻².

3. DERIVATION OF CURTIS MATRIX ELEMENTS

The net upward flux of thermal radiation at wavenumber ν is given by

$$F_{\nu}(\tau) = 2\pi \int_0^1 d\mu \int_0^{\infty} J_{\nu}^{+}(t) e^{-|t-\tau|/\mu} d|t-\tau| - 2\pi \int_0^1 d\mu \int_0^{\infty} J_{\nu}^{-}(t) e^{-|t-\tau|/\mu} d|t-\tau|, \quad (3.1)$$

(Goody, 1964, p. 49)

where μ is the absolute value of the cosine of the zenith angle, and the optical depth is defined as

$$\tau = \tau(p) = \frac{1}{g} \int_0^p k_{\nu}(p'') c(p'') dp'', \quad (3.2)$$

where k_{ν} is the absorption coefficient and c the concentration of the absorbing gas. The source function J_{ν} is assumed to be isotropic in the upper (+) and lower (-) hemispheres. Changing the variable of integration to pressure yields

$$F_{\nu}(p) = 2\pi \int_0^1 d\mu \int_p^{p_s} J_{\nu}^{+}(p) \exp \left\{ - \int_p^{p'} k_{\nu} \frac{cdp''}{g\mu} \right\} k_{\nu} \frac{c}{g} dp' + 2\pi \int_0^1 d\mu \int_{p_s}^{\infty} J_{\nu}^{+}(p_s) \exp \left\{ - \int_p^{p'} k_{\nu} \frac{cdp''}{g\mu} \right\} k_{\nu} \frac{c}{g} dp' - 2\pi \int_0^1 d\mu \int_p^0 J_{\nu}^{-}(p') \exp \left\{ - \int_{p'}^p k_{\nu} \frac{cdp''}{g\mu} \right\} k_{\nu} \frac{c}{g} dp', \quad (3.3)$$

where p_s is the surface pressure. The second term, arising from the surface, can be integrated over p' directly.

Consider the flux of radiation integrated over a frequency interval $\Delta\nu$ sufficiently small that J_ν^\pm can be approximated by mean values over frequency J^\pm . Then changing the order of integration gives

$$\begin{aligned} \int_{\Delta\nu} d\nu F_\nu(p) = & 2\pi \int_p^{p_s} dp' J^+(p') \int_0^1 d\mu \int_{\Delta\nu} d\nu \exp \left\{ - \int_p^{p'} k_\nu \frac{cdp''}{g\mu} \right\} k_\nu \frac{c}{g} \\ & + 2\pi J^+(p_s) \int_0^1 d\mu \mu \int_{\Delta\nu} d\nu \exp \left\{ - \int_p^{p_s} k_\nu \frac{cdp''}{g\mu} \right\} \\ & - 2\pi \int_p^0 dp' J^-(p') \int_0^1 d\mu \int_{\Delta\nu} d\nu \exp \left\{ - \int_{p'}^p k_\nu \frac{cdp''}{g\mu} \right\} k_\nu \frac{c}{g} \quad (3.4) \end{aligned}$$

Note that each of the integrals in the exponents are non-negative.

Define the flux equivalent width

$$e(p, p') = \int_0^1 d\mu \mu \int_{\Delta\nu} d\nu [1 - T_\nu(p, p', \mu)]$$

where

$$T_\nu(p, p', \mu) = \exp \left\{ - \int_p^{p'} k_\nu c / g\mu dp'' \right\} \quad (3.5)$$

Function $e(p, p')$ is simply one minus the transmission function between p and p' weighted by μ integrated over one hemisphere. Clearly $e(p, p) = 0$, and

$$\frac{\partial e}{\partial p'} = \int_0^1 d\mu \int_{\Delta\nu} d\nu \exp \left\{ - \left| \int_p^{p'} k_{\nu} c \frac{dp'}{g\mu} \right| \right\} \frac{k_{\nu} c}{g} \times \begin{cases} 1, & p' > p \\ -1, & p' < p \end{cases} \quad (3.6)$$

Changing variables from p to e , the expression for the integrated flux becomes

$$\begin{aligned} \int_{\Delta\nu} d\nu F_{\nu}(p) = 2\pi \int_0^{e(p, p_s)} J^+(p') de(p, p') + 2\pi J^+(p_s) \left[\frac{\Delta\nu}{2} - e(p, p_s) \right] \\ + 2\pi \int_{e(p, 0)}^0 J^-(p') de(p, p') \quad (3.7) \end{aligned}$$

Define the upward flux

$$F^+(p) = 2\pi \int_0^{e(p, p_s)} J^+(p') de(p, p') + 2\pi J^+(p_s) \left[\frac{\Delta\nu}{2} - e(p, p_s) \right]$$

and the downward flux

$$F^-(p) = 2\pi \int_0^{e(p, 0)} J^-(p') de(p, p') \quad (3.8)$$

Assume the source function is isotropic over all directions $J^+ = J^-$, and define a nondimensional "temperature" $\theta = J/\bar{B}$, where \bar{B} is the Planck function in $\Delta\nu$ evaluated at some reference temperature. We have

$$\begin{aligned} \int_{\Delta\nu} d\nu F_{\nu}(p) = F^+(p) - F^-(p) = 2\pi\bar{B} \int_0^{e(p, p_s)} \theta(p') de(p, p') \\ + 2\pi\bar{B} \left[\frac{\Delta\nu}{2} - e(p, p_s) \right] - 2\pi\bar{B} \int_0^{e(p, 0)} \theta(p') de(p, p') \quad (3.9) \end{aligned}$$

3.1 Vertical grid

Divide the atmosphere into N layers (see Fig. 1). The pressure at ϕ_1 is zero, and ϕ_{N+1} is at the ground. Pressure and temperature evaluated at levels within layers are denoted with carets. To simplify notation, let

$$e(p_1, p_j) = e_{1j} \quad e(p_1, \hat{p}_j) = \hat{e}_{1j} \quad e(p_1, p_s) = e_{1,N+1}.$$

For the simplest case θ is treated as a constant throughout each layer. Then

$$\begin{aligned} F_1^+ &= 2\pi\bar{B} \sum_{j=1}^N (e_{1,j+1} - e_{1j}) \hat{\theta}_j + 2\pi\bar{B} \left(\frac{\Delta v}{2} - e_{1,N+1} \right) \hat{\theta}_{N+1} \\ &= 2\pi\bar{B} \sum_{j=1}^{N+1} (e_{1,j+1} - e_{1j}) \hat{\theta}_j \quad \text{where } e_{1,N+2} \equiv \Delta v/2, \end{aligned} \quad (3.10)$$

and

$$F_1^- = -2\pi\bar{B} \sum_{j=1}^{i-1} (e_{1,j+1} - e_{1j}) \hat{\theta}_j.$$

The convergence of radiant flux in layer i is given by

$$\begin{aligned} \Delta F_i &= F_{i+1}^+ - F_{i+1}^- - F_i^+ + F_i^- \\ &= 2\pi\bar{B} \sum_{j=1}^{N+1} (e_{i+1,j+1} - e_{i+1,j} - e_{i,j+1} + e_{i,j}) \hat{\theta}_j. \end{aligned} \quad (3.11)$$

The change in temperature in layer i resulting from the convergence of radiant flux is

$$q_1 = \sum_{j=1}^{N+1} R_{1j} \hat{\theta}_j \text{ where } R_{1j} = \frac{2\pi \bar{B}g}{\Delta p_1 c_p} (e_{i+1,j+1} - e_{i+1,j} - e_{i,j+1} + e_{ij}) , \quad (3.12)$$

where Δp_1 is the pressure difference across layer 1.

A more complex integration scheme results if we assume that temperature varies linearly in log pressure from layer midpoint to layer midpoint:

$$\theta(\phi) = \gamma_1 \hat{\theta}_1 + (1 - \gamma_1) \hat{\theta}_{i-1}, \quad \hat{\phi}_{i-1} \leq \phi \leq \hat{\phi}_1$$

where $\gamma_1 = (\ln \phi - \ln \hat{\phi}_{i-1}) / (\ln \hat{\phi}_1 - \ln \hat{\phi}_{i-1})$. Above $\hat{\phi}_1$ the atmosphere is assumed to be isothermal at θ_1 , and below $\hat{\phi}_N$ the temperature is extrapolated from level $\hat{\phi}_{N-1}$ to a temperature discontinuity at the surface. The expression for the Curtis matrix elements becomes

$$R_{1j} = \frac{2\pi \bar{B}g}{\Delta p_1 c_p} (e'_{i+1,j+1} - e'_{i+1,j} - e'_{i,j+1} + e'_{ij}) ,$$

where $e'_{ij} = e_{ij} + c_{ij}$ and correction term c_{ij} is given by

$$c_{ij} = \begin{cases} 0, & \text{if } j = 1 \\ \hat{e}_{ij} - e_{ij} - \int_{\hat{e}_{i,j-1}}^{\hat{e}_{ij}} \gamma_j de, & 2 \leq j \leq N-1 \\ e_{i,N+1} - e_{i,N} - \int_{e_{i,N-1}}^{e_{i,N+1}} \gamma_N de, & j = N \\ 0, & \text{if } j = N+1 \text{ or } N+2 \end{cases} \quad (3.13)$$

The integrals in (3.13) have been performed numerically using the trapezoidal rule.

3.2 The overlap of spectral lines

Below about 30 km pressure-broadened line widths may be sufficiently large to cause significant overlapping of individual rotational lines. The entire band (550-800 cm^{-1}) is divided into ten equal spectral subintervals, and spectral lines are assumed to be randomly distributed in position within each subinterval. Eqn. (3.5) can be written

$$e(p_i, p_j) = \Delta\nu \left[\frac{1}{2} - \int_0^1 d\mu \mu \frac{1}{\Delta\nu} \sum_{k=1}^{10} \delta\nu_k \bar{T}_k \right] , \quad (3.14)$$

where $\bar{T}_k = \frac{1}{\delta\nu_k} \int_{\delta\nu_k} d\nu T_\nu$ is the mean transmission function over subinterval k.

The assumption of random distribution implies (Goody, 1964, p. 154)

$$\bar{T}_k = \exp \left(- \frac{1}{\delta\nu_k} \sum W_\ell \right) , \quad (3.15)$$

where W_ℓ is the equivalent line width of the ℓ^{th} line in subinterval k.

3.3 Line shapes

Through much of the region of interest the shape of the individual spectral lines is given by the Voigt profile, a convolution of the simpler Doppler and Lorentz shapes. The equivalent width of a Voigt line can be adequately approximated with a formula that interpolates between the equivalent line widths given by the Doppler and Lorentz profiles (Rodgers and Williams, 1974)

$$W(p_i, p_j, \mu) = (W_L^2 + W_D^2 - W_L^2 W_D^2 / S^2 u^2)^{1/2} \quad (3.16)$$

where

$$W_L = W_L[u/\mu, S(\tilde{T}_{1j}), \alpha_L(\tilde{p}_{1j}, \tilde{T}_{1j})]$$

and

$$W_D = W_D[u/\mu, S(\tilde{T}_{1j})] \quad .$$

W_L is the Lorentz equivalent line width given by an accurate approximation due to Goldman (1968). W_D is the Doppler equivalent line width given by a series approximation or a large Su approximation (Goody, 1964, p. 136).

$S(\tilde{T}_{1j})$ is the line strength evaluated at the Curtis-Godson temperature \tilde{T}_{1j} between pressure levels p_1 and p_j and u is the amount of absorber (CO_2) between levels p_1 and p_j . The Lorentz line width $\alpha_L(\tilde{p}_{1j}, \tilde{T}_{1j})$ is given by

$$\alpha_L(\tilde{p}_{1j}, \tilde{T}_{1j}) = 0.08 \text{ cm}^{-1} \left(\frac{\tilde{p}_{1j}}{p_0} \right) \left(\frac{300 \text{ K}}{\tilde{T}_{1j}} \right)^{1/2}$$

where the Curtis-Godson parameters are

$$\tilde{T}_{1j} = \frac{1}{p_1 - p_j} \int_{p_j}^{p_1} T(p) \, dp \quad \text{and} \quad \tilde{p}_{1j} = \frac{1}{2} (p_1 + p_j) \quad .$$

For a given temperature profile the Curtis-Godson temperature was computed for every pair of levels and stored for future reference as a two-dimensional array.

3.4 Variation of line strength with temperature

There are nearly 17,000 CO_2 lines tabulated between 550 and 800 cm^{-1} on the AFGL Atmospheric Absorption Line Parameter Compilation ("McClatchey tape"). Rather than sum over all these lines in (3.15), the lines in each subinterval were divided into groups according to line strengths, and a mean line strength was used to represent the strength of all the lines in

that group. Working with one subinterval at a time, the strength of each line was computed at seven temperatures between 150 K and 296 K with the formula

$$S(T) = \frac{S(T_S)Q_v(T_S)Q_r(T_S)}{Q_v(T)Q_r(T)} \exp \left\{ \frac{1.439E''(T-T_S)}{TT_S} \right\}, \quad (3.17)$$

where Q_v and Q_r are the vibrational and rotational partition functions interpolated from the table in McClatchey et al. (1973). T_S is a reference temperature, and E'' is the energy of the lower state of the vibrational-rotational transition expressed in cm^{-1} . S is expressed in the units $\text{cm}^{-1} (\text{cm STP})^{-1}$. At each temperature the lines were sorted into 50 groups representing equal increments of $\log_{10} S$. The resulting histograms of line strengths for one subinterval are shown in Fig. 2. The histograms were represented approximately by dividing the lines into two categories: a set of temperature-independent strong lines, and a larger set of weaker lines which follow the distribution function

$$n_1(x) = m [x - x_{01}]^2 \text{ for } x < x_{01}, \quad (3.18)$$

where $x = \log_{10} S$. Here $x_{01} = \log_{10} S_{01}$ where S_{01} is the strength corresponding to the empty bin of lines just stronger than the strongest temperature dependent lines in subinterval 1. Factor m was determined for each temperature in each subinterval, but was found to be adequately represented by a universal function

$$m(T) = aT + b, \quad a = 0.001615, \quad b = -0.2066, \quad (3.19)$$

where a and b are evaluated by regression. The mean strength in each group is given by

$$\bar{S}_j = \exp(2.303 \bar{x}_j) \text{ where } \bar{x}_j = \int dx n(x) x / \int dx n(x) , \quad (3.20a)$$

and the number of lines having this strength is

$$F_j = \int dx n(x) . \quad (3.20b)$$

The number and mean strength of the strong lines in each subinterval was evaluated directly. These strong lines, together with weak lines evaluated from (3.20) were used to compute mean transmission functions, flux equivalent widths, and finally, the Curtis matrix elements for a standard temperature profile.

3.5 Non-LTE source function (this section follows Houghton, 1977, p. 58ff)

Above about 65 km the frequency of collisions is insufficient to maintain a population of excited vibrational states consistent with the local kinetic temperature. The source function for a particular transition is no longer identical to the Planck function, but is given by

$$J_\nu = \frac{\bar{L}_\nu + \epsilon B_\nu}{1 + \epsilon} , \quad (3.21)$$

where $\epsilon = C/A$ is the ratio of the rate of collisional deactivation to the rate of spontaneous radiative deactivation of the excited state. Assuming J_ν and B_ν are constant over the frequency band, and using the equation of radiative transfer to eliminate the mean radiance \bar{L}_ν gives

$$J = B + q c_p / 4\pi \epsilon S_b , \quad (3.22)$$

where c is the concentration of the absorber, and S_b and q are the local band strength and radiative heating of the particular transition. Eqn. (3.12) must be modified to

$$q_i = \sum_{j=1}^{N+1} R_{ij} (J_j / B) \quad (3.23)$$

Using matrix notation, (3.22) and (3.23) can be combined into

$$q = R^{nLTE} \theta \quad \text{where } R^{nLTE} = (I - RE)^{-1} R \quad (3.24)$$

and

$$E_{ij} = [c_p / 4\pi B c \epsilon_i S_{b,i}] \delta_{ij} \quad .$$

$S_{b,i}$ and ϵ_i are the band strength and LTE factor evaluated in the i^{th} layer.

In the present calculation all lines are considered part of the fundamental band with Einstein coefficient $A = 0.74 s^{-1}$ and collision frequency $C = 0.67 \times 10^5 s^{-1}$ at standard pressure and an assumed mesopause temperature of 180K. The band strength $S_{b,i}$ is evaluated as

$$S_{b,i} = \sum_{k=1}^{10} \left(\sum_j F_{kj} \bar{S}_{kj} + F_k^s S_k^s \right) \quad , \quad (3.25)$$

where F_{kj} and \bar{S}_{kj} are given by (3.20) and F^s and S^s refer to the strong lines in subinterval k .

3.6 Temperature dependence of transmission function

Since the flux equivalent widths depend on the temperature of the intervening material, the Curtis matrix elements are functions of the

atmospheric temperature profile. The flux equivalent width is related to the transmission function, so we have adapted the interpolation method of Fels and Schwarzkopf (1981).

Let $k_\nu(p'') = \sum_{\ell} f_{\ell}[\nu, T(p''), p'']$ where f_{ℓ} is the absorption coefficient of the ℓ^{th} line in $\delta\nu$. Let $T(p'') = T_0(p'') + \delta T(p'')$ where $T_0(p'')$ is the standard temperature profile. To first order in δT

$$\begin{aligned} T_{\nu}(p, p', \mu) &= \exp \left\{ - \int_p^{p'} \sum_{\ell} \left[f_{\ell}(\nu, T_0(p''), p'') + \left. \frac{\partial f_{\ell}}{\partial T} \right|_{T_0} \delta T(p'') \right] \frac{cdp''}{g\mu} \right\} \\ &= \exp \left\{ - \int_p^{p'} \sum_{\ell} f_{\ell}(\nu, T_0(p''), p'') \frac{cdp''}{g\mu} \right\} \exp \left\{ - \int_p^{p'} \sum_{\ell} \left. \frac{\partial f_{\ell}}{\partial T} \right|_{T_0} \delta T(p'') \frac{cdp''}{g\mu} \right\} . \end{aligned}$$

Assume $\delta T(p'')$ is small and expand the second exponent

$$T_{\nu}(p, p', \mu) = T_{\nu}^0(p, p', \mu) \left[1 - \int_p^{p'} \sum_{\ell} \left. \frac{\partial f_{\ell}}{\partial T} \right|_{T_0} \delta T(p'') \frac{cdp''}{g\mu} \right] ,$$

where

$$T_{\nu}^0(p, p', \mu) = \exp \left\{ - \int_p^{p'} \sum_{\ell} f_{\ell}(\nu, T_0(p''), p'') \frac{cdp''}{g\mu} \right\} .$$

Equ. (3.5) becomes

$$\begin{aligned} e(p, p') &= \int_0^1 d\mu \int_{\Delta\nu} \left[1 - T_{\nu}^0(p, p', \mu) \right] \\ &\quad - \int_p^{p'} \frac{cdp''}{g} \int_0^1 d\mu \int_{\Delta\nu} T_{\nu}^0(p, p', \mu) \sum_{\ell} \left. \frac{\partial f_{\ell}}{\partial T} \right|_{T_0} \delta T(p'') . \end{aligned}$$

Now define $e^0(p, p') = \int_0^1 d\mu \mu \int_{\Delta\mu} dv [1 - T_v^0(p, p', \mu)]$ and

$$\frac{\delta e(p, p', p'')}{\delta T} = - \frac{c}{g} \int_0^1 d\mu \int_{\Delta\mu} dv T^0(p, p', \mu) \sum_k \frac{\partial f_k}{\partial T} [v, T(p''), p''] \Big|_{T_0},$$

so

$$e(p, p') = e^0(p, p') + \int_p^{p'} dp'' \frac{\delta e(p, p', p'')}{\delta T} \delta T(p'') .$$

Consider the functional derivative $\delta e / \delta T$. Suppose it can be factored (and in several approximations it can) as

$$\frac{\delta e(p, p', p'')}{\delta T} = F(p, p') G(p''),$$

then

$$\delta e(p, p') = \int_p^{p'} dp'' \frac{\delta e(p, p', p'')}{\delta T} \delta T(p'') = \int_p^{p'} dp'' F(p, p') G(p'') \delta T(p'')$$

$$= F(p, p') \int_p^{p'} dp'' G(p'') \delta T(p'') .$$

If we let $e^+(p, p')$ represent the flux equivalent width for a temperature profile everywhere warmer than T_0 by a constant amount $\bar{\Delta}$, and similarly for $e^-(p, p')$, then from a Taylor series we have for any constant temperature deviation δT

$$\begin{aligned} e(p, p') &= e^0(p, p') + \frac{e^+(p, p') - e^-(p, p')}{2\bar{\Delta}} \delta T \\ &+ \frac{\frac{1}{2}[e^+(p, p') - 2e^0(p, p') + e^-(p, p')]}{(\bar{\Delta})^2} (\delta T)^2 . \end{aligned} \quad (3.26)$$

Now define the weighted temperature deviation

$$\Delta(p, p') = \frac{\int_p^{p'} dp'' G(p'') \delta T(p'')}{\int_p^{p'} dp'' G(p'')} \quad (3.27)$$

We will assume that for variable temperature deviations (3.26) holds with $\Delta(p, p')$ replacing δT . The best choice for weighting function $G(p)$ and for $\bar{\Delta}$ have been determined by Fels and Schwarzkopf as

$$G(p) = p^{0.2} [1 + (p/p_0)^{0.8}] \quad (3.28)$$

where $p_0 = 30$ mb, and $\bar{\Delta} = 25K$.

Therefore the procedure adopted has been to compute a matrix of e , R , and R^{nLTE} for each of three temperature profiles. Each R^{nLTE} is then "unraveled" using (3.12) to obtain a matrix e^{nLTE} . This can be accomplished because there are "boundary conditions" $e_{11} = 0$ and $e_{1,N+2} = \Delta\nu/2$. Beginning with the lowest level, successive layers of e_{ij} can be computed from (3.12), first on the lower side of the diagonal ($j > i$), and then on the upper ($j < i$). These three arrays are stored as data for the dynamical model. When the radiative algorithm is called, $e^{nLTE}(T)$ is computed by (3.26) for the appropriate temperature profile, then $R^{nLTE}(T)$ is computed from (3.12).

The Curtis matrix for the 10μ band of CO_2 is computed using the same procedure. The only difference is that the 10μ band was not divided into subintervals to compute the mean transmission functions of (3.14).

4. INFRARED COOLING BY OZONE

The version of the algorithm currently incorporated in the dynamical model uses a cool-to-space plus cool-to-ground approximation utilizing the band absorption formula of Ramanathan (1976). In this section we shall develop a more general method from which the above approximation can be easily derived.

Several empirical and semiempirical formulae (e.g., Walshaw, 1957, Aida, 1975), have been developed for the integrated absorption A of the 9.6μ band of ozone where

$$A = \int_{\Delta\nu} d\nu A_\nu = \int_{\Delta\nu} d\nu (1 - T_\nu) \quad (4.1)$$

These formulae may be used to compute flux equivalent width and Curtis matrix elements from (3.5) and (3.12). The procedure outlined in section 3 may be greatly simplified. IR cooling (or heating) by ozone is significant only below about 60 km, so non-LTE effects can be ignored. The contribution of the 9.6μ band to the total cooling is small enough that variations of the transmission function with temperature profile can be neglected. However since the vertical distribution of ozone varies, in general, with latitude and time, a new Curtis matrix must be computed for every ozone profile.

Use of the Curtis matrix method facilitates the identification of the various components of the radiative interactions. The radiative cooling is given by

$$q_1 = \sum_{j=1}^{N+1} R_{1j} \theta_j$$

where $N+1$ is the ground. If the atmosphere is isothermal, and the ground is at the same temperature, then radiation exchanged with the atmosphere and ground provides no net cooling and there is only cooling to space. So

$$q_1^{\text{CTS}} = \sum_{j=1}^{N+1} R_{1j} \theta_1 = \theta_1 \sum_{j=1}^{N+1} R_{1j} = \theta_1 C_1^s \quad (4.2)$$

If the ground is at a different temperature, there can be cooling to both space and ground. Then

$$\begin{aligned} q_1^{\text{CTS}} + q_1^{\text{CTG}} &= \sum_{j=1}^{N+1} R_{1j} \theta_j = \theta_1 \sum_{j=1}^N R_{1j} + R_{1,N+1} \theta_{N+1} \\ &= \theta_1 C_1^s + (\theta_{N+1} - \theta_1) C_1^g \end{aligned} \quad (4.3)$$

where $C_1^g = R_{1,N+1}$.

A further approximation can be introduced for θ . The "temperature" is given by

$$\theta = \frac{e^u - 1}{e^{uT_0/T} - 1} \quad \text{where } u = hv/kT_0 = 4.988 \text{ for } T_0 = 300\text{K}.$$

Hence, $e^u - 1 \approx e^u$ and for $T \leq T_0$ we have $e^{uT_0/T} - 1 \approx e^{uT_0/T}$,

and $\theta \approx e^{-uT_0/T} e^{-u}$.

$$\begin{aligned} \text{So } q_1^{\text{CTS}} + q_1^{\text{CTG}} &\approx e^u C_1^s e^{-uT_0/T} + e^u C_1^g (e^{-uT_0/T_g} - e^{-uT_0/T_1}) \\ &\approx 146.6 C_1^s e^{-1496/T} + 146.6 C_1^g (e^{-1496/T_g} - e^{-1496/T_1}), \end{aligned} \quad (4.4)$$

where T is given in Kelvins. The radiative cooling of the 9.6μ band for the U.S. Standard ozone and temperature profiles computed with the Curtis matrix method employing Aida's parameterization is compared with the cool-to-space plus cool-to-ground approximation in Fig. 3.

5. PROGRAMS AND SUBROUTINES

5.1 Routines used to obtain line strength distributions

The 16899 lines attributed to carbon dioxide between 550 and 800 cm^{-1} have been sorted into eight sets according to the energy of the lower state of the transition. Each set is written on an element file of the fileset CO2SET. The description of each element set is as follows ("normal" refers to $^{12}\text{C}^{16}\text{O}_2$):

File name	Lower level	Number of lines	Type of CO_2
FU626	0	153	"normal"
HO626	ν_2	546	"normal"
HT626	$\nu_1 \approx 2\nu_2$	159	"normal"
HD626	$2\nu_2$	691	"normal"
HF626	$3\nu_2$	564	"normal"
HC626	$\geq 4\nu_2$	3805	"normal"
FUIISO	0	1269	isotopic
HTISO	$\geq \nu_2$	9712	isotopic

Program BINS sorts all of these lines into bins according to strength without regard to band of origin. The entire band is divided into ten equal subintervals and the strength of each line is computed at seven different temperatures. The width of each bin is 0.25 in units of $\log S$, where S is expressed in $(\text{cm}^{-1}) (\text{cm STP})^{-1}$. An example of the resulting histograms is shown in Fig. 2. If the McClatchey tape is used as the input for program BINS, then DO loop 12 should be replaced with $\text{IB} = 8$, the two IF statements after statement number 98 should be activated, and LMAX must be set equal to or greater than 16899.

Program FIT fits each of the seventy histograms produced by BINS with a second order polynomial, according to (3.18). The only free parameter is the coefficient m .

Graphs of m vs. T for each of the ten subintervals are shown in Fig. 4. It is sufficiently accurate to assume a universal function $m(T) = aT + b$ for all subintervals. Program REGRESS provides the least square values for a and b .

Spectral lines associated with the fundamental have no vibrational temperature dependence and are usually strong. They can be identified on the histograms because they rise above the generally parabolic shape of the line strength distribution profile (see Fig. 2). Program RESID subtracts the fitted line strength distribution from the actual distribution, and computes the number of strong lines and their mean strength in each bin, and then consolidates each four bins to give the number of strong lines and their mean strength in each strength decade. From this information one characteristic strong line number and strength is chosen for each subinterval, and specified in program CURTEL in arrays FS and SS.

5.2 Routines used to obtain Curtis matrices

Program CURTEL computes the Curtis matrix elements for the 15μ band of carbon dioxide. Wave numbers are expressed in units of the Doppler width at 300K.

The computation of the non-LTE Curtis matrix begins with the evaluation of the total band strength by summing all the mean strengths weighted by the number of lines at that strength over all strength decades and subintervals. VMULFM and LINVIF are IMSL routines for matrix multiplication and inversion.

Subroutine FLUXEL calculates the equivalent band widths. First the pressure levels (PHI) and layer midpoints (PHIM) are computed. The Curtis-Godson temperature (TCG) between any two levels is computed and stored. The correction terms c_{ij} [see (3.13)] are evaluated by the trapezoidal rule with the interval divided into M1 panels for the diagonal terms, M panels for the nondiagonal terms, and $3*M/2$ panels for the terms at the top of the atmosphere.

In subroutine EFUNC there is an opportunity to use the integrated absorption parameterizations of Pollack et al. (1981) or Rodgers and Walshaw (1966) instead of the line strength distributions that have been derived. If the integrated absorption parameterizations are not chosen, the Curtis-Godson temperature and pressure are evaluated and the line strength distribution parameter [m in (3.18), called A2 in this subroutine] is calculated.

Subroutine W performs the summation in (3.14), evaluating the mean transmission function for the entire band. The line strength distribution parameter is used to compute the number of lines and mean strengths in nine decades of line strength for each of the ten subintervals. This synthetic line distribution is cut off at the low strength end so that the total number of lines in each subinterval is the same for the synthetic and real distributions.

After the set of non-LTE Curtis matrices have been produced for the three temperature profiles and for both 10μ and 15μ bands, program EQWD inverts (3.12) to obtain a set of equivalent band widths. The set of six arrays of $e^{n_{LTE}}$ are used by subroutine CURT to compute the temperature-dependent Curtis matrix elements.

There are, unfortunately, two different conventions for the labeling of the layers of the model. In the dynamical model (JRH) the lowest stratospheric layer is denoted $j = 1$, and the highest layer is $j = \text{JML}$ (i.e., 16). In the radiation algorithm (CBL) the highest layer is denoted $j' = 1$, the lowest stratospheric layer is $j' = \text{JML}$, the troposphere is $j' = \text{JML} + 1$, and the ground is $j' = \text{JML} + 2$ (i.e., 18). Comment cards in the listing noting the shift of conventions will hopefully prevent confusion.

5.3 Routines used for ozone and oxygen

Program OZPRO produces the arrays needed to compute the radiative heating and cooling due to ozone, given the ozone mixing ratio and temperature profiles. Subroutine COLOZ computes the specific density and column abundance of ozone. Subroutine CMOZ96 computes a Curtis matrix for the 9.6μ band, and from it cool-to-space and cool-to-ground coefficients are derived. Two different parameterizations (Aida, 1975 and Walshaw, 1957) of the integrated absorption of the 9.6μ band are offered in function OZ96A.

Subroutine OXY computes the density and column abundance of molecular oxygen.

Function OZUV computes the solar energy absorbed per molecule of O_2 from the parameterization by Strobel (1978).

Function OZUV computes the derivative of the fraction of total solar flux absorbed with respect to ozone column abundance from the parameterization of Lacis and Hansen (1974).

5.4 Routines used in the dynamical model

Subroutine HEAT computes the diabatic heating term for the dynamical equations. The earth-sun geometry factors and the set of Curtis matrices are computed once per day. The static stability is computed once, during the initial time step of the run. The zonal mean heating is introduced into the dynamical equations gradually with a growth time (GTIME) of two days.

Subroutine RADEQU computes the global radiative equilibrium temperature profile (i.e., the temperature profile such that the radiative heating integrated over the entire globe is zero at each level).

Function DELT computes the individual terms of the net radiative heating. The parameterization of Lacis and Hansen (1974) is used to compute the radiative heating due to absorption of solar visible and UV radiation by ozone. Values of the various parameters describing the shortwave planetary albedo are taken from the program of Ramanathan (private communication). The parameterization of Strobel (1978) is used for the absorption of solar radiation by molecular oxygen.

5.5 Changes in main program WAVE2

The flux equivalent width arrays are read in first for 10μ , then 15μ , in the order minus, zero, plus. KAP is the Newtonian cooling coefficient (day^{-1}) derived from the present algorithm, assuming a U.S. Standard temperature profile and the ozone profile (based on Krueger and Minzner, 1976) currently used in the dynamical model. The coefficients have been computed for a "boxcar" perturbation of vertical scale 25 km. TZO is on input the U.S. Standard atmosphere temperature profile. Array PRS contains the Fourier

decomposition of the observed \bar{u} at the lower boundary. Q_{OG} and Q_O are $Q_3(\text{ground})$ and Q_3 of (2.1). Q_{OZSG} and Q_{OZS} are $n_3(\text{ground})$ and $n_3(p_k)$ of (2.2), CTS is $-146.6C^B$ and CG is $146.6C^B$, defined in (4.4).

Subroutine HEAT is called every IRCT time steps; IRAD is used as a counter.

Tropospheric and ground temperatures are computed with simple equations that fit the climatological temperatures to within a few degrees.

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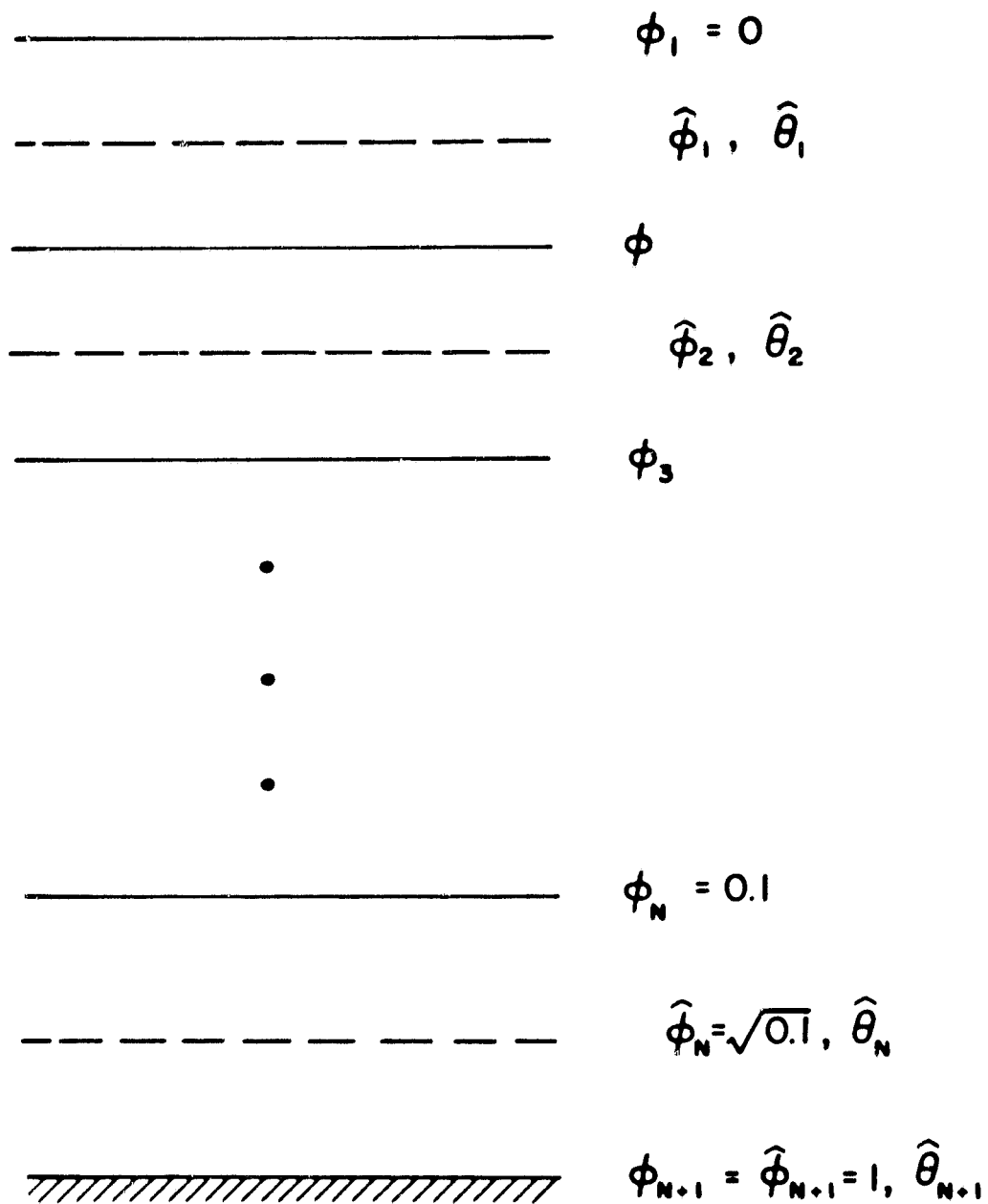


Figure 1: Vertical structure of atmospheric model.

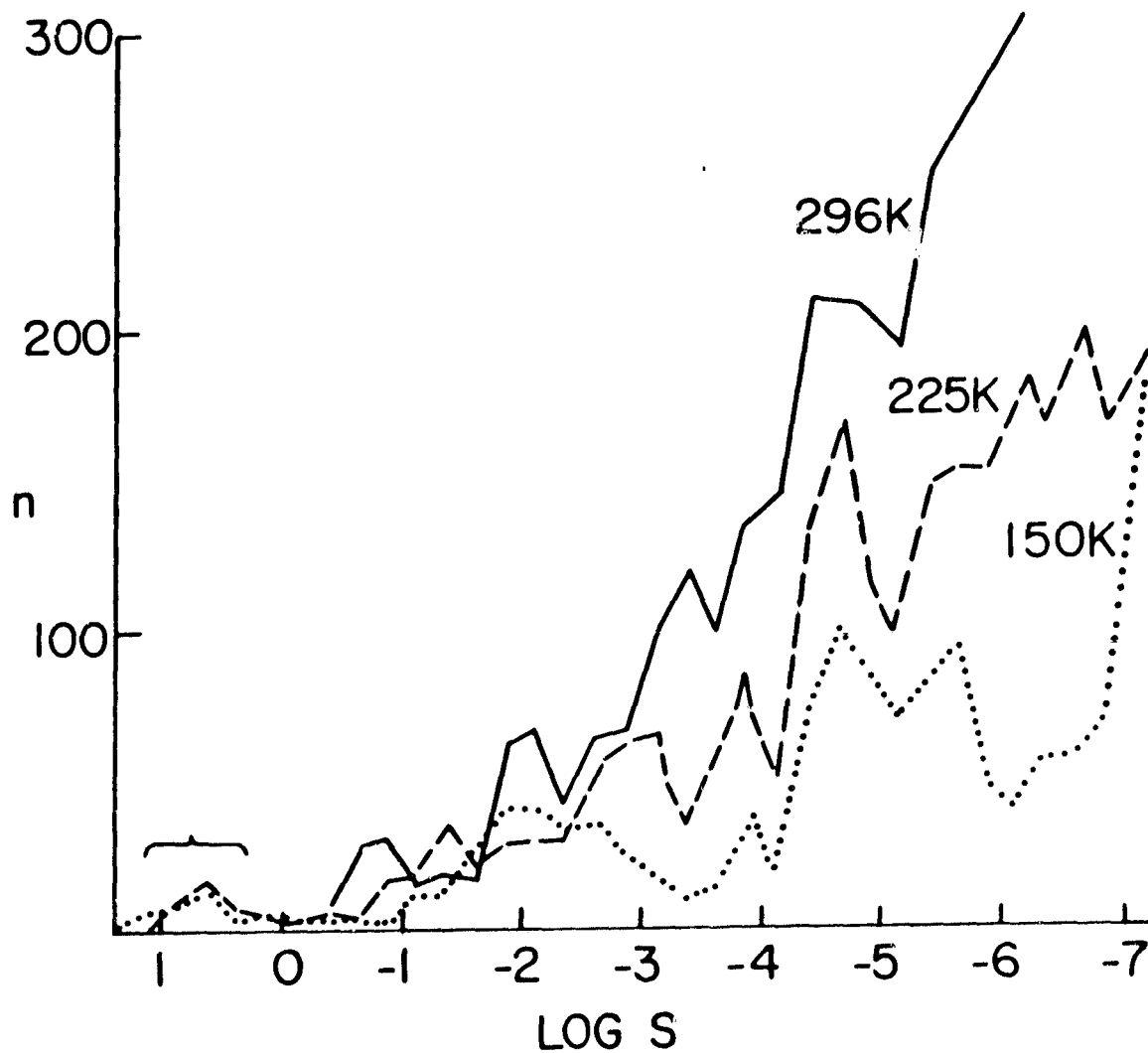


Figure 2: Histograms of line strength distribution at three different temperatures for subinterval $650 \leq \nu \leq 675 \text{ cm}^{-1}$. Increments of $\log S$ are 0.25. Strength S expressed in units of $\text{cm}^{-1}(\text{cm STP})^{-1}$. Lines under bracket are "strong lines", whose temperature dependence has been neglected.

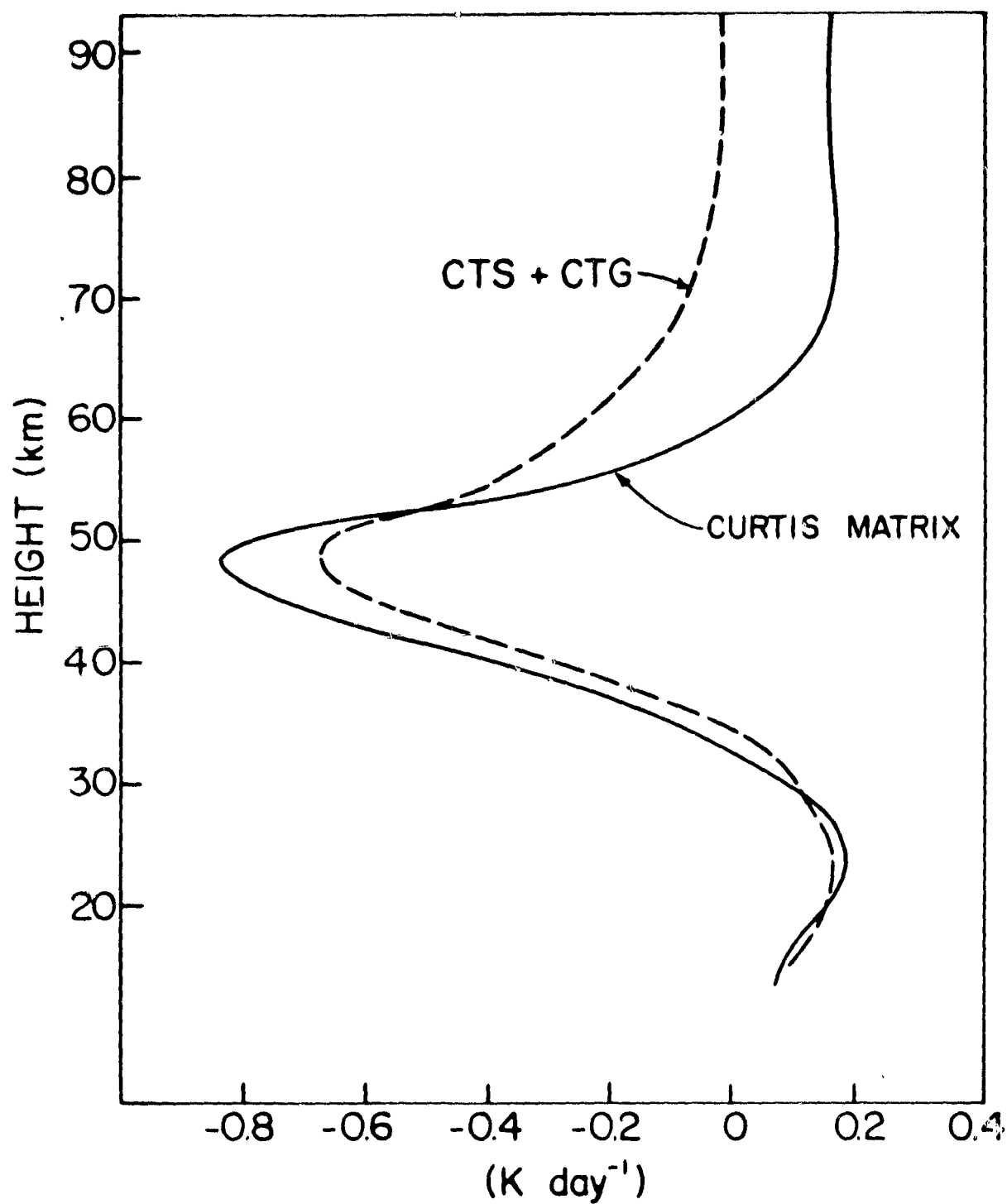


Figure 3: Comparison of cool-to-space plus cool-to-ground approximation (broken line) with Curtis matrix method (solid line) for radiative cooling by the 9.6μ band of ozone. Temperature and ozone profiles are from the 1976 U.S. Standard Atmosphere.

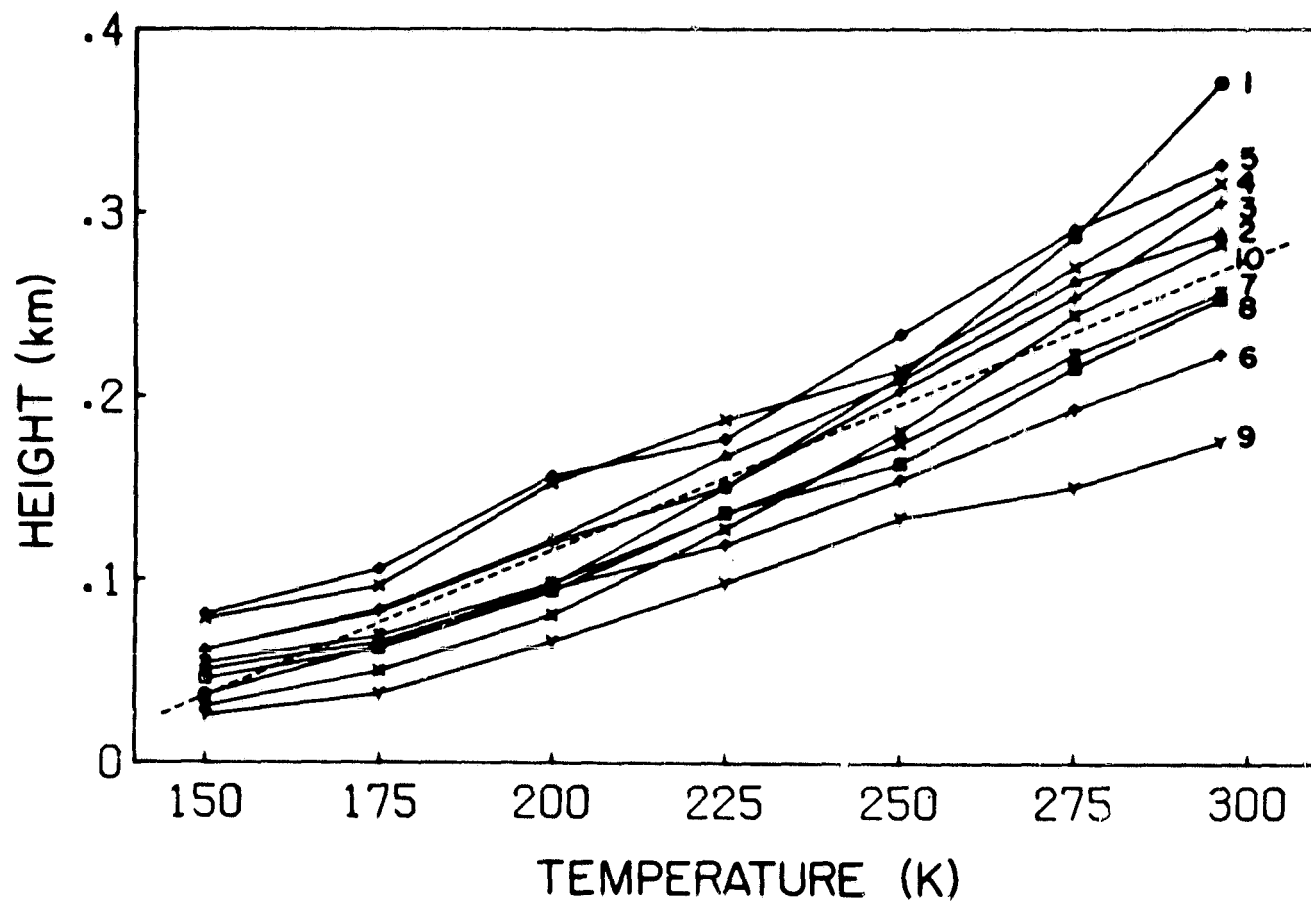


Figure 4: Weak line distribution parameter m vs temperature. Solid lines connect values for the same subinterval, identified on right (e.g., subinterval 1 is $550 \leq \nu \leq 575 \text{ cm}^{-1}$). Broken line is best fit for all subintervals.

APPENDIX

FORTRAN Code

```

      PROGRAM BINS(INPUT,OUTPUT,TAPE1,TAPE2,TAPE3,TAPE4,TAPE5
      1,TAPE6,TAPE7,TAPE8,TAPE9)
C   SUBROUTINE TO SORT LINES ACCORDING TO STRENGTH
C   (TAPES 1 THROUGH 8 ARE ELEMENT FILES OF CO2SET)
C
      DIMENSION N(10,7,50)
      DIMENSION TP(7),QV(7)
      INTEGER US
      DATA TP/150.,175.,200.,225.,250.,275.,296./
      DATA QV/1.004,1.0095,1.0142,1.0327,1.0502,1.0719,1.0931/
      TS=296.      QVS=1.0931      RJ=1.
      OSTP=2.6869E+19
C   FROM OUTPUT APRAYS
      DO 11 IS=1,10
      DO 9 IT=1,7
      DO 3 I=1,50
      N(IS,IT,I)=0
      3 CONTINUE
      9 CONTINUE
      11 CONTINUE
C   SORT EACH BAND
      DO 12 IB=1,8
      IS=1
      UPLIM=575.
      READ(IB,15) LMAX
      LL=0
      99 LL=LL+1
      IF(LL.GT. LMAX) GO TO 101
      98 READ(IU,5) FREQ,S,ALPHA,ELS,US,LS,ROT,DATE,ISO,MOLC
      IF(FREQ.LT. 550.) GO TO 98
C   IF(MOLC.NE. 2) GO TO 98
      IF(FREQ.GT. UPLIM) GO TO 101
      97 CONTINUE
      DO 7 IT=1,7
      I=TP(IT)
      EF=EXP(1.430*ELS*(T-TS)/(T+TS))
      ST=5*EF*QVS/QV(IT)*(1/TS)**(-RJ)
C   EXPRESS STRENGTH IN UNITS (CM ATM)-1 CM-1
      ST=ST*OSTP
      SLT=ALOG10(ST)
C   CHOOSE CORRECT BIN
      IG=50
      DO 31 I=1,49
      P=I-50      RL=1.25+9/4.
      IF(SLT.GT. RL) IG=50-I
      31 CONTINUE
C   COUNT LINES IN EACH BIN AND STUKE PARAMETERS
      N(IS,IT,IG)=N(IS,IT,IG)+1
      7 CONTINUE
      GO TO 99
      101 CONTINUE
      IF(IB.NE. 8) GO TO 13
      DL=UPLIM-25.
      WRITE(9,35) UL,UPLIM
      DO 8 IT=1,7
      WRITE(4,45) TP(IT)
      DO 4 I=1,50
      NN=N(IS,IT,I)
      IF(NN.EQ. 0) GO TO 4
      WRITE(9,25) I,NN
      4 CONTINUE
      IP=51
      WRITE(9,25) IP
      8 CONTINUE
      13 UPLIM=UPLIM+25.
      IS=IS+1
      IF(UPLIM.GT. 800.) GO TO 12
      GO TO 97
      12 CONTINUE
      25 FORMAT(10X,2I6,E15.3)
      5  FORMAT(F10.3,E10.3,F5.3,F10.3,1I2,3X,1I2,A7,A4,A4,2X,1I)
      15 FORMAT(10X,I5)
      35 FORMAT(* SPECTRAL INTERVAL*,F5.0,* TO*,F5.0)
      45 FORMAT(10X,* TEMPERATURE*,F5.0)
      END

```

```

      PROGRAM FIT(INPUT,OUTPUT,TAPE1,TAPE2)
C     SUBROUTINE TO FIT HISTOGRAMS WITH PARABOLAS
C     (TAPE1 IS OUTPUT OF PROGRAM BINS)
C
      DIMENSION N(51),LC(10)
      DATA LC/16,9,8,3,0,1,5,9,11,16/
C     LC IS BIN NUMBER OF WEAKEST EMPTY BIN IN EACH SUBINTERVAL
      DO 4 IS=1,10
      READ(1,45) DUM
      DO 3 IT=1,7
      READ(1,5) T
      P=1.
      TOT=0.
      DO 1 I=1,51
      N(I)=0
1     CONTINUE
      DO 2 I=1,51
      READ(1,15) I1,N(I1)
      IF(I1 .EQ. 51) GO TO 6
      TN=N(I1)
      TOT=TOT+TN
      IF(TN .LT. P) GO TO 2
      NC=I1
      P=TN
2     CONTINUE
6     IF(TOT .EQ. 0.) GO TO 3
      IF(NC .GT. 40) NC=40
      IC1=LC(IS)
      NP=NC-IC1+1
      S1=S2=0.
      DO 7 I=1,NP
      S1=S1+N(IC1+I-1)*(I-1)*(I-1)
      S2=S2+(I-1)**4
7     CONTINUE
      RM=S1/S2
      WRITE(2,25) IS,LC(IS),T,NC,TOT,RM
3     CONTINUE
      WRITE(2,35)
4     CONTINUE
5     FORMAT(23X,F5.0)
25    FORMAT(2I5,F10.1,I5,F10.1,2F10.5,E10.4)
35    FORMAT(/)
45    FORMAT(10A6)
15    FORMAT(10X,2I6,E15.5)
      END

```



```

      PROGRAM REGRESS(INPUT,OUTPUT,TAPE1)
C     SUBROUTINE TO COMPUTE BEST VALUES FOR UNIVERSAL FIT
C     (TAPE1 IS OUTPUT OF PROGRAM FIT)
C
      REAL M,M
      SM=ST2-SMT-ST=0.
      N=0.
      DO 2 I=1,10
      DO 3 IT=1,7
      READ(1,5) T,M
      PRINT 9, T,M
      N=N+1.
      ST=ST+T
      SMT=SMT+M*T
      SM=SM+M
      ST2=ST2+T*T
3     CONTINUE
      READ(1,15) DUM
      READ(1,15) DUM
2     CONTINUE
      D=N*ST2-ST*ST
      A=(N*SMT-SM*ST)/D
      B=(SM*ST2-SMT*ST)/D
      PRINT 25, A,B
5     FORMAT(10X,F10.1,15X,F10.5)
15    FORMAT(10A6)
25    FORMAT(' A=*,E15.5,10X,* B=*,E15.5)
      END

```

```

      PROGRAM RESID(INPUT,OUTPUT,TAPE1)
C     SUBROUTINE TO COMPUTE RESIDUAL STRONG LINES IN 15 MICRON BAND
C     (TAPE1 IS OUTPUT OF PROGRAM BINS)
C
      DIMENSION N(10,7,51),XZ(10),A(7),SR(7),SS(7),R(7)
      DATA XZ/16 9.,8.,3.,0.,1.,5.,9.,1.,16./
      DO 4 IS=1,10
      READ(1,65) DUM
      DO 3 IT=1,7
      READ(1,5) T
      DO 1 I=1,51
      N(IS,IT,I)=0.
2     CONTINUE
      A(IT)=0.001589*T-0.2017
      DO 2 I=1,51
      READ(1,15) II,N(IS,IT,II)
      IF(II.EQ. 51) GO TO 3
2     CONTINUE
3     CONTINUE
4     CONTINUE
      DO 6 IS=1,10
      PRINT 55
C     RESIDUES FOR STRONGEST DECADE OF LINE STRENGTHS
      DO 14 IT=1,7
      RR=N(IS,IT,1)
      SS(IT)=0
      R(IT)=0.
      IF(1.LE. XZ(IS)) GO TO 16
      R(IT)=RR-A(IT)*(1.-XZ(IS))*2
16     SR(IT)=R(IT)
      E=-0.25*(1.-5.5)
      S=10**E
      IF(KR.EQ. 0.) S=0.
      SS(IT)=S
14     CONTINUE
      PRINT 35, (SR(IT),IT=1,7)
      PRINT 45, (SS(IT),IT=1,7)
C     RESIDUES FOR SEVEN MORE DECADES OF LINE STRENGTHS
      DO 9 K=1,7
      DO 12 J=1,7
      SR(J)=SS(J)=0.
12     CONTINUE
      DO 8 J=1,4
      I=4*(K-1)+J+1
      RI=I
      DO 11 IT=1,7
      RR=N(IS,IT,I)
      P(IT)=0.
      IF(RI.LE. XZ(IS)) GO TO 13
      R(IT)=RR-A(IT)*(RI-XZ(IS))*2
13     SR(IT)=SR(IT)+P(IT)
      E=-0.25*(RI-5.5)
      S=10**E
      SS(IT)=SS(IT)+S*R(IT)
11     CONTINUE
      PRINT 25, I, (R(II),II=1,7)
8     CONTINUE
      DO 17 IT=1,7
      IF(SR(IT).EQ. 0.) GO TO 17
      SS(IT)=SS(IT)/SR(IT)
17     CONTINUE
C     MEAN FOR FOUR BINS (ONE DECADE)
      PRINT 35, (SR(IT),IT=1,7)
      PRINT 45, (SS(IT),IT=1,7)
9     CONTINUE
6     CONTINUE
25     FORMAT(15,7F10.1)
35     FORMAT(7F15.1)
45     FORMAT(7E15.3)
15     FORMAT(10X,2I6,E15.5)
55     FORMAT(1H1)
5     FORMAT(23X,F5.0)
65     FORMAT(10A6)
      END

```

```

PROGRAM CURTEL(INPUT,OUTPUT,TAPE6,TAPE7)
COMMON//S(10,50),F(10,10),DELNU,ALPHDZ,E(41,40),E(41,40)
S,PHI(40),PHIH(40),TCG(20,20),TT(21)
DIMENSION T(21),R(41,40),Q(20),CTS(20),THETA(20)
DIMENSION H(21,20),U(20,20),UI(20,20),WK(20),RR(20,21),QO(20)
DIMENSION FS(10),SS(10)
REAL NUZ,KB,MASS
DATA N,M,M1,L,DELTANU/17,4,16,3,250./
DATA FS/2*0.,20.,5.,22.,7.,4*0./
DATA SS/2*0.,0.05,2.,5.5,4.2,4*0./
DATA T/187.1,186.4,193.2,201.6,210.4,221.0,234.7,248.4,263.9,
1 3470.7,258.0,243.2,229.6,224.2,219.6,216.7,231.4,288.2,3*0./
PI=3.1415926
SP1=SQRT(PI)
HP=6.6237E-27
AVG=6.0225E+23
MASS=44./AVG
C=2.9979E+10
KB=1.3805E-16
CP=1.E+07
GRAV=981.
TZ=300.
C PARAMETERS FOR 15 MICRON BAND
NUZ=667.379
D=EXP(HP*C*NUZ/(KB*TZ))-1.
BBAZ=2.*HP*(C*NUZ)**3/(C*D)
CONST=2.*PI*BBAZ*GRAV/CP
ALPHDZ=NUZ*SQRT(2.*KB*TZ/MASS)/C
DELNU=DELTANU
C
N1=N+1
DO 2 I=1,10
F(I,1)=FS(I) S S(I,1)=SS(I)/ALPHDZ
DO 27 J=2,10
F(I,J)=S(I,J)=0.
27 CONTINUE
2 CONTINUE
DO 26 I=1,N1
TT(I)=T(I)
26 CONTINUE
PRINT 101, N,M,M1,L,DELNU
CALL FLUXEL(N,M,M1,L)
DO 13 I=1,10
DO 14 J=1,10
S(I,J)=S(I,J)*ALPHDZ
14 CONTINUE
WRITE(6,103)((F(I,J),S(I,J)),J=1,10)
13 CONTINUE
C
C COMPUTE LTE CURTIS MATRIX
DO 20 I=1,N
CTS(I)=0.
DP=1.013E+06*(PHI(I+1)-PHI(I))
DO 21 J=1,N1
R(J,I)= E1(J+1,I+1)+E1(J,I)-E1(J+1,I)-E1(J,I+1)
R(J,I)=(CONST/DP)*R(J,I)*86400.
CTS(I)=CTS(I)+R(J,I)
21 CONTINUE
20 CONTINUE
C
WRITE(6,100)((R(J,I),J=1,N1),I=1,N)
WRITE(6,100)(T(I),I=1,N1)
WRITE(6,102)
WRITE(6,100) (CTS(I),I=1,N)
WRITE(6,102)
C
C COMPUTE TEMPERATURE THETA = B/B(REF)
VZ=HP*C*NUZ/(KB*TZ)
DO 16 J=1,N1
V=HP*C*NUZ/(KB*T(J))
THETA(J)=(EXP(VZ)-1.)/(EXP(V)-1.)
16 CONTINUE
C

```

```

C   COMPUTE HEATING RATES
      DO 10 I=1,N
        Q(I)=0.
        DO 9 J=1,N1
          Q(I)=Q(I)+R(J,I)*THETA(J)
        9 CONTINUE
        WRITE(6,105) Q(I)
      10 CONTINUE

C
C   COMPUTE NON-LTE CURTIS MATRIX
      CM=0.00033*44./28.966
      SBAND=0.
C   COMPUTE TOTAL BAND STRENGTH FROM DISTRIBUTION
      DO 23 I=1,10
        DO 24 J=1,10
          SBAND=SBAND+F(I,J)*S(I,J)
        24 CONTINUE
      23 CONTINUE
      SBAND=SBAND*22414./44.

C
C   PERFORM MATRIX ALGEBRA
      DO 6 J=1,N1
        DO 7 K=1,N
          H(J,K)=0.
C   EPS = COLLISION RATE/SPONTANEOUS EMISSION RATE
          EPS=0.74*PHI(H(K))/1.5E-05
          IF(J .EQ. K) H(J,K)=CP/(4.*PI*CM*SBAND*EPS*86400.)
        7 CONTINUE
      6 CONTINUE
      CALL VMULFH(P,H,N1,N,N,41,21,U,20,IER)
      PRINT 101, IER
      DO 8 I=1,N
        DO 11 J=1,N
          U(I,J)=-U(I,J)
          IF(I .EQ. J) U(I,J)=U(I,J)+1.
        11 CONTINUE
      8 CONTINUE
      CALL LINVIF(U,N,20,U1,0,WK,IER)
      PRINT 101, IER
      DO 17 I=1,N
        DO 18 K=1,N1
          SUM=0.
          DO 19 J=1,N
            SUM=SUM+U1(I,J)*R(K,J)
          19 CONTINUE
          RR(K,I)=SUM
        18 CONTINUE
      17 CONTINUE
      WRITE(6,102)

C   NON LTE CURTIS MATRIX
      WRITE(6,100)((RR(J,I),J=1,N1),I=1,N)
      WRITE(7,100)((RR(J,I),J=1,N1),I=1,N)

C   COMPUTE HEATING RATE FOR NON-LTE CURTIS MATRIX
      DO 3 I=1,N
        CQ(I)=0.
        DO 4 J=1,N1
          CQ(I)=CQ(I)+RR(J,I)*THETA(J)
        4 CONTINUE
        WRITE(6,105) CQ(I)
      3 CONTINUE
      DO 12 I=1,N
        CTS(I)=0.
        DO 15 J=1,N1
          CTS(I)=CTS(I)+RR(J,I)
        15 CONTINUE
      12 CONTINUE
      WRITE(6,100) (CTS(I),I=1,N)
105  FORMAT(E10.3)
103  FORMAT(10E12.4)
102  FORMAT(/)
100  FORMAT(9F8.3/9F8.3)
101  FORMAT(4I3,F6.1)
      STOP
      END

```

```

      SUBROUTINE FLUXELIN,M,M1,L)
C      SUBROUTINE FOR CALCULATING FLUX EQUIVALENT BAND WIDTHS
      CCH10N//S(10,10),F(10,10),DELNU,ALPHOZ,E(41,40),E(41,40)
      ,PHI(40),PHIH(40),TCG(20,20),T(21)
      DIMENSION C(41,40)
      N1=M+1
      N2=M+2
      DEL= DELNU/2.
      IF(L .EQ. 2) DEL=.45.

C
      GO TO I=1,M
      PHI(I)= PFUNC(I,M)
10     PHIH(I)=PHFUNC(I,M)
      PHI(N1)= 1.
      PHIH(N1)=SQRT(0.1)
      PHIH(N1)=1.

C
C      COMPUTE CURTIS-GOOSON TEMPERATURE FOR EACH PAIR OF LEVELS
      DO 22 I=1,M1
      DO 23 J=I,M1
      TCG(I,J)=T(I)
      IF(I .EQ. J) GO TO 23
      SUM=0.
      JM1=J-1
      DO 26 L=I,JM1
      SUM=SUM+(T(L)+T(L+1))*(PHIH(L+1)-PHIH(L))
26     CONTINUE
      TCG(I,J)=0.5*SUM/(PHIH(J)-PHIH(I))
      TCG(J,I)=TCG(I,J)
23     CONTINUE
      PRINT 101,(TCG(I,J),J=1,M1)
22     CONTINUE

C
      DO 70 I=1,M1
      DO 50 J=1,M
      IF (J.EQ.1) GO TO 46
      IF (J.EQ.1) GO TO 30
      IF (J.GE.N) GO TO 48

C
      E(J,1)= EFUNC(PHI(J),PHI(I),L)
      COR=0.

C      COMPUTE NONDIAGONAL TERM
      DO 15 K=1,M
      EK1=K-1
      EK=K
      PHIK= ((M-EK1)/M)*PHIH(J-1)+(EK1/M)*PHIH(J)
      PHIK1= ((M-EK)/M)*PHIH(J-1)+(EK/M)*PHIH(J)
15     COR= COR+(GFUNC(PHIK,PHIH(J),PHIH(J-1))+GFUNC(PHIK1,PHIH(J)
      ,PHIH(J-1)))*(EFUNC(PHIK1,PHI(I),L)-EFUNC(PHIK,PHI(I),L))

C
      C(J,I)= EFUNC(PHIH(J),PHI(I),L)-E(J,I)-COR/2.
      GO TO 50

C
C      COMPUTE DIAGONAL TERM
30     CORDN=0.
      CORUP=0.
      E(I,I)=0.

C
      DO 35 K=1,M1
      EK1=K-1
      EK=K
      PHIK= ((M1-EK1)/M1)*PHIH(I-1)+(EK1/M1)*PHI(I)
      PHIK1= ((M1-EK)/M1)*PHIH(I-1)+(EK/M1)*PHI(I)
      CORDN= CORDN+(GFUNC(PHIK,PHIH(I),PHIH(I-1))+GFUNC(PHIK1,PHIH(I),
      PHIH(I-1)))*(EFUNC(PHIK,PHI(I),L)-EFUNC(PHIK1,PHI(I),L))
      PHIK= ((M1-EK1)/M1)*PHI(I)+(EK1/M1)*PHIH(I)
      PHIK1= ((M1-EK)/M1)*PHI(I)+(EK/M1)*PHIH(I)
35     CORUP= CORUP+(GFUNC(PHIK,PHIH(I),PHIH(I-1))+GFUNC(PHIK1,
      PHIH(I),PHIH(I-1)))*(EFUNC(PHIK1,PHI(I),L)-EFUNC(PHIK,PHI(I),L))

C
      C(I,I)= EFUNC(PHIH(I),PHI(I),L)+(CORDN-CORUP)/2.

C
      GO TO 50

C

```

```

C   COMPUTE NONDIAGONAL TERM AT TOP
46  C(I,I)=0.
    E(I,I)= EFUNC(0.,PHI(I),L)
    GO TO 50
C
48  E(N,I)= EFUNC(PHI(N),PHI(I),L)
    E(N1,I)= EFUNC(PHI(N1),PHI(I),L)
    COR=0.
    M2= 3*M/2
    DO 49 K=1,M2
      EK1=K-1
      EK=K
      PHIK= ((M2-EK1)/M2)*PHI(N-1)+(EK1/M2)*PHI(N1)
      PHIK1= ((M2-EK)/M2)*PHI(N-1)+(EK/M2)*PHI(N1)
49  COR = COR+(GFUNC(PHIK,PHI(N),PHI(N-1))+GFUNC(PHIK1,PHI(N)
    1,PHI(N-1)))+(EFUNC(PHIK1,PHI(1),L)-EFUNC(PHIK,PHI(1),L))
    C(N,I)= E(N1,I)-E(N,I)-COR/2.
C
50  E1(J,I)=E(J,I)+C(J,I)
C
    E1(N1,I)=EFUNC(PHI(N1),PHI(I),L)
70  E1(N2,I)= DEL
C
    RETURN
    END

```

```

FUNCTION PFUNC(J,N)
  RJN=J-N
  PFUNC=0.1*EXP(RJN*5./7.)
  IF(J.EQ. 1) PFUNC=0.
  RETURN
END

```

```

FUNCTION PHFUNC(J,N)
C   THIS SUBROUTINE COMPUTES PRESSURE HALF LEVELS
  PHFUNC=SQRT(PFUNC(J,N)*PFUNC(J+1,N))
  IF(J.EQ. 1) PHFUNC=0.5*PFUNC(2,N)
  RETURN
C   END
FUNCTION GFUNC(A,B,C)
C   PLANCK FUNCTION LINEARLY INTERPOLATED IN LOG PRESSURE.
  GFUNC= ALOG(A/C)/ALOG(B/C)
  RETURN
END

```

```

FUNCTION EFUNC(A,B,L)
COMMON//S(10,10),F(10,10),DELNU,ALPHDZ,E1(41,40),E(41,40)
S,PHI(40),PHIH(40),TCG(20,20),IT(21)
DIMENSION Y(2)
DATA Y/0.2113,0.7667/
RLEV(X)=7./5.*ALOG(10.*X)+17.
IF(L.EQ. 1) GO TO 10
IF(L.EQ. 2) GO TO 20
IF(L.EQ. 3) GO TO 30

C
C   POLLACK 15 MICRON BAND,SLP=1000MB,C=.00033.
10 EFUNC=26.55*ALOG(1.+1.1386*((1926.43*ABS(A-B))*0.566)
1* ((A+B)*0.323))
GO TO 80

C
C   RODGERS-WALSHAW 15 MICRON BAND,SLP=1000MB,C=.00033 BY VOLUME
C   EVALUATED AT 200K
20 EFUNC=0.
U=ABS(A-B)
IF(U.EQ. 0.) GO TO 80
P=(A+B)/2.
TRAN=0.
DO 21 J=1,2
UP=U/Y(J)
TRAN=TRAN+0.5*Y(J)*EXP(-300.3*UP/SQRT(1.+541.3*UP/P))
21 CONTINUE
EFUNC=170.*(0.5-TRAN)
GO TO 80

C
C   DOPPLER-LORENTZ LINE
30 GRAV=981.
EFUNC=0.
C   COMPUTE PATH LENGTH
U=(B-A)+1.013E+06*0.00033*22414./(GRAV*28.966)
U=ABS(U)
IF (U.EQ. 0.) GO TO 80
TZ=300.
C   CUKTIS-GOUSON APPROXIMATION
IA=1
IF(A.EQ. 0.) GO TO 2
IA=RLEV(A)
EP=RLEV(A)-IA
IF(EP.GE. 0.5) IA=IA+1
2 IB=1
IF(B.EQ. 0.) GO TO 3
IB=RLEV(B)
EP=RLEV(B)-IB
IF(EP.GE. 0.5) IB=IB+1
3 IF(IA.GT. 18) IA=18
IF(IB.GT. 18) IB=18
T=TCG(IA,IB)
C   COMPUTE WEAK LINE DISTRIBUTION PARAMETER
SL=0.001615
BINT=-0.2066
A2=SL*T+BINT
P=(A+B)/2.
PZ=1.
AL=0.08/ALPHDZ*P/PZ*TZ/T
C   TWO POINT GAUSSIAN QUADRATURE OVER ANGLES
TRAN=0.
DO 22 J=1,2
TRAN=TRAN+0.5*Y(J)*W(U/Y(J),AL,A2)
22 CONTINUE
EFUNC=DELNU*(0.5-TRAN)

C
80 CONTINUE
RETURN
END

```

```

      FUNCTION W(U,A,A2)
C     THIS SUBROUTINE COMPUTES MEAN TRANSMISSION FUNCTION
      COMMON/7S(10,10),F(10,10),DELNU,ALPHDZ,E1(41,40),E(41,40)
      DIMENSION XZ(10),N(10)
      DATA XZ/16.,9.,5.,3.,0.,1.,5.,9.,11.,16./
      DATA N/683,1407,1676,2640,3724,2360,1912,1270,734,493/
C     N IS TOTAL NUMBER OF LINES IN EACH SUBINTERVAL
C     XZ IS ONE LESS THAN NUMBER OF FIRST OCCUPIED BIN
      PI=3.1415926
      SPI=SQRT(PI)
      W=0.
      DNU=DELNU/10.
C     TEN SPECTRAL INTERVALS
      DO 6 I=1,10
        XM=XZ(I)+(3.*N(I)/A2)**(1./3.)
C     XM NORMALIZES CONTINUOUS DISTRIBUTION FUNCTION TO N
        U1=(XM-5.5)/4.
        U2=(XZ(I)-9.5)/4.
        DO 8 J=2,10
          P=J-4
          IF(R.GT. U1) GO TO 8
          IF(R.LT. U2) GO TO 8
          VU=4.*R+9.5
          IF(XM.LT. VU) VU=XM
          VL=4.*R+5.5
          IF(XZ(I).GT. VL) VL=XZ(I)
C
C     COMPUTE NUMBER OF LINES IN STRENGTH DECADE
          F(I,J)=A2/3.*((VU-XZ(I))**3-(VL-XZ(I))**3)
C     ESTIMATE MEAN STRENGTH IN DECADE
          XMN=0.25*A2*((VU-XZ(I))**4-(VL-XZ(I))**4)/F(I,J)+XZ(I)
          S(I,J)=10**(-0.25*(XMN-5.5))
          S(I,J)=S(I,J)/ALPHDZ
        8 CONTINUE
        7 CONTINUE
        WW=0.
C     SUM OVER LINE STRENGTH DECADES
        DO 4 J=1,10
          IF(S(I,J).EQ. 0.) GO TO 4
          SU=S(I,J)*U
C     GOLDMAN APPROXIMATION FOR LORENTZ LINE EQUIVALENT WIDTH
          WL=SU/(1.+(0.25*SU/A)**1.25)**0.4
          IF(SU.GT. 4.) GO TO 2
C     SERIES APPROXIMATION FOR DOPPLER LINE EQUIVALENT WIDTH
          WD=WD1(SU/SPI)
          GO TO 3
C     LARGE SU APPROXIMATION FOR DOPPLER LINE EQUIVALENT WIDTH
        2 U2=SQRT(ALOG(SU/SPI))
          WD=2.*(U2+0.2896/U2-0.1335/U2**3+0.0070/U2**5)
C     INTERPOLATION FOR VOIGT LINE
        3 WV=SQRT(WL*WL+WD*WD-WL*WL*WD*WD/(SU*SU))
          WW=WW+F(I,J)*WV
        4 CONTINUE
          W1=EXP(-WW*ALPHDZ/DNU)
          W=W+W1*DNU
        6 CONTINUE
          W=W/DELNU
          RETURN
      END

```



```

      FUNCTION WD1(W)
C   THIS SUBROUTINE COMPUTES DOPPLER LINE EQUIVALENT WIDTH FROM SERIES
      PI=3.14159265
      SPI=SQRT(PI)
      NF=1
      SUM=0.
      DO 2 N=1,6
      RN=N
      NF=NF*N
      R=NF
      SUM=SUM+(-W)*N/(R*SQRT(RN))
2  CONTINUE
      WD1=-SPI*SUM
      IF(WD1 .LE. 0.) WD1=1.
      RETURN
      END

```

```

      PROGRAM EQUW(INPUT,OUTPUT,TAPES,TAPES)
C     THIS SUBROUTINE COMPUTES EQUIVALENT BAND WIDTH FROM CURTIS MATRICES
      DIMENSION R(10,17),PHI(10),E(17,10),CST(2),DN(2)
      REAL KB,NUZ(2)
      PI=3.1415926
      HP=6.6237E-27
      C=2.9979E+10
      KB=1.3805E-16
      CP=1.E+07
      GRAV=981.
      TZ=300.
      DN(2)=250.      & DN(1)=258.
      NUZ(2)=667.379
      NUZ(1)=979.
C     FOR 10 AND 15 MICRON BANDS
      DO 1 M=1,2
        U=EXP(HP*C*NUZ(M)/(KB*TZ))-1.
        BBAR=2.*HP*(C*NUZ(M))*3/(C*U)
        CONST=2.*PI*BBAR*GRAV/CP
        CST(M)=CONST
      1 CONTINUE
      N=17      & N1=N+1      & N2=N+2      & NM=N-1
C     CBL VERTICAL GRID
      DO 10 I=1,N
        PHI(I)=PFUNC(I,N)
      10 CONTINUE
      PHI(N1)=1.
C     CALCULATE FUNCTION E FOR SIX CURTIS MATRICES:
C     R=(10),RO(10),R+(10),R-(15),RO(15),R+(15)
      DO 2 L=1,6
        CONST=CST(1)      & DELNU=DN(1)
        IF(L.LT.4) GO TO 3
        CONST=CST(2)      & DELNU=DN(2)
      3 CONTINUE
      READ(5,25)((R(J,I),J=1,N1),I=1,N)
      DO 20 I=1,N
        DP=1.013E+06*(PHI(I+1)-PHI(I))
        DO 21 J=1,N1
          R(J,I)=DP*R(J,I)/(CONST*86400.)
      21 CONTINUE
      20 CONTINUE
      DO 9 I=1,N1
        E(I,I)=0.
      9 CONTINUE
C
      DO 4 I=1,N1
        E(N2,1)=DELNU/2.
      4 CONTINUE
C
      DO 7 JJ=1,N1
        J=N2-JJ
        JM=J-1
        DO 8 II=1,JM
          I=J-II
          E(J,I)=R(J,I)-E(J+1,I+1)+E(J+1,I)+E(J,I+1)
      8 CONTINUE
      7 CONTINUE
      DO 11 KK=1,N
        K=KK-1
        NK=N-K
        DO 12 L=1,NK
          I=N1-L
          J=1-K
          E(J,I+1)=-R(J,I)+E(J+1,I+1)+E(J,I)-E(J+1,I)
      12 CONTINUE
      11 CONTINUE
      WRITE(6,5)((E(J,I),J=1,N2),I=1,N1)
      DO 30 I=1,N
        DP=1.013E+06*(PHI(I+1)-PHI(I))
        DO 31 J=1,N1
          R(J,I)=E(J+1,I+1)+E(J,I)-E(J+1,I)-E(J,I+1)
          R(J,I)=R(J,I)*CONST*86400./DP
      31 CONTINUE
      30 CONTINUE

```

```
C  PRINT A TO CHECK PROGRAM
    PRINT 25, ((R(J,I),J=1,10),I=1,17)
  2  CONTINUE
  5  FORMAT(J(5E20.13/),4E20.13)
 25  FORMAT(9F8.3)
    END
```

	SUBROUTINE CURT(TT,KN,TST,K)	CURT
C	SUBROUTINE TO COMPUTE TEMPERATURE-DEPENDENT CURTIS MATRICES	CURT
C	(TT IS ACTUAL TEMPERATURE PROFILE, TST IS STANDARD TEMPERATURE	
C	PROFILE)	
C	COMMON/CURTIS/PII(18),PIIH(18),G1(17),G2(18),E(6,19,10)	CURT
	6,CUMAT(36,17,18)	CURT
	DIMENSION T(18),TS(18),EL(2,19,18),TT(18),TST(18),R(2,18,17),	CURT
	SB(18,18),U(18,18),DT(18)	CURT
	REAL KB,NUZ(2)	CURT
	KNL=KN-1 8 KP=KN+1 8 K2=(N+2	CURT
C	CONVERT TO CGL VERTICAL COORDINATE CONVENTION	CURT
	DO 2 K=1,KNL	FEB3
	TS(K)=TST(KN-K)	FEB3
	T(K)=TT(KN-K)	FEB3
2	CONTINUE	FEB3
	DO 8 K=KN,KP	FEB3
	TS(K)=TST(K)	FEB3
	T(K)=TT(K)	FEB3
8	CONTINUE	FEB3
C	COMPUTE TEMPERATURE DEVIATION FROM STANDARD	CURT
	DO 1 I=1,KP	CURT
	DT(I)=T(I)-TS(I)	CURT
	U(I,I)=DT(I)	CURT
	B(I,I)=1.	CURT
1	CONTINUE	CURT
C	COMPUTE WEIGHTED DEVIATION	CURT
	DO 3 I=1,KN	CURT
	IP=I+1	CURT
	DO 4 J=IP,KP	CURT
	U(I,J)=DT(I)*G1(I)+DT(J)*G2(J)	CURT
	B(I,J)=G1(I)+G2(J)	CURT
	IF(J.EQ. KP) D(I,J)=U(I,J)+U.5*(DT(KP)-DT(KN))*G1(KN)	CURT
	L1=I+1 8 L2=J-1	CURT
	IF(L1.GT. L2) GO TO 7	CURT
	DO 6 L=L1,L2	CURT
	D(I,J)=U(I,J)+DT(L)*(G1(L)+G2(L))	CURT
	B(I,J)=B(I,J)+G1(L)+G2(L)	CURT
6	CONTINUE	CURT
7	CONTINUE	CURT
4	CONTINUE	CURT
3	CONTINUE	CURT
C	NORMALIZE DEVIATION	CURT
	DO 9 I=1,KN	CURT
	IP=I+1	CURT
	DO 11 J=IP,KP	CURT
	D(J,I)=D(I,J)=D(I,J)/B(I,J)	CURT
11	CONTINUE	CURT
9	CONTINUE	CURT
C	COMPUTE TEMPERATURE-DEPENDENT EQUIVALENT (BAND) WIDTHS	CURT
	DO 12 M=1,2	CURT
	DO 13 I=1,KP	CURT
	DO 14 J=1,K2	CURT
	EE(M,J,I)=E(3*M-1,J,I)+E(3*M,J,I)-E(3*M-2,J,I)*D(J,I)/50.+	CURT
	50.5*(E(3*M,J,I)+L(3*M-2,J,I)-2.*E(3*M-1,J,I))*D(J,I)+	CURT
	SD(J,I)/625.	CURT
14	CONTINUE	CURT
13	CONTINUE	CURT
12	CONTINUE	CURT
C	COMPUTE CURTIS MATRICES	CURT
	PI=3.141592654	CURT
	HP=6.6237E-27	CURT
	C=2.9979E+10	CURT
	KB=1.3805E-16	CURT
	CP=1.E+07	CURT
	GRAV=981.	CURT
	TZ=300.	CURT
	NUZ(2)=667.379	CURT
	NUZ(1)=980.	CURT
	DO 19 M=1,2	CURT
	UD=EXP(HP*C*NUZ(M)/(KB*TZ))-1.	CURT
	BBAR=2.*HP*(C*NUZ(M))+3/(C*UD)	CURT
	CUNST=2.*PI*BBAR*GRAV/CP	CURT

CURT
CURT
CURT
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PROGRAM OZPRO(INPUT,OUTPUT,TAPE1,TAPE2)
COMMON//E1(41,40),E(41,40)
S,PHI(40),PHIH(40),TCG(20,20),Q(18),QZS(18)
DIMENSION QUG(18),QO(18,17),QOZSG(18),QOZS(18,17),CTS(18,17),
CTG(18,17)
DIMENSION W(18),CS(17),CG(17)
JM=19 6 KN=17
JML=JM-1 5 KNL=KN-1
KNP=KN+1
C COMPUTE OZONE PARAMETERS FOR EACH LATITUDE ZONE
DO 2 J=1,1
READ(1,5)(W(K),K=1,KNP)
READ(1,15)(T(K),K=1,KNP)
C W IS MASS MIXING RATIO IN G/G AT PRESSURE HALF LEVELS
C NUMBERED DOWNWARD FROM TOP OF ATMOSPHERE. T IS TEMPERATURE AT SAME L LEVELS
CALL GLOBE(KN,W,T,CS,CG)
QUG(J)=Q(KNP)
QOZSG(J)=QZS(KNP)
QU(J,KN)=Q(KN)
QOZS(J,KN)=QZS(KN)
CTS(J,KN)=CS(KN)
CTG(J,KN)=CG(KN)
DO 3 K=1,KNL
QO(J,K)=Q(KN-K)
QOZS(J,K)=QZS(KN-K)
CTS(J,K)=CS(KN-K)
CTG(J,K)=CG(KN-K)
3 C/CONTINUE
2 CONTINUE
WRITE(2,195)(QUG(J),J=1,JML)
WRITE(2,195)(QO(J,KN),J=1,JML)
WRITE(2,195)((QO(J,K),J=1,JML),K=1,KNL)
WRITE(2,195)(QOZSG(J),J=1,JML)
WRITE(2,195)(QOZS(J,KN),J=1,JML)
WRITE(2,195)((QOZS(J,K),J=1,JML),K=1,KNL)
WRITE(2,195)(CTS(J,KN),J=1,JML)
WRITE(2,195)((CTS(J,K),J=1,JML),K=1,KNL)
WRITE(2,195)(CTG(J,KN),J=1,JML)
WRITE(2,195)((CTG(J,K),J=1,JML),K=1,KNL)
195 FORMAT(9E10.3)
5 FORMAT(EY.2)
15 FORMAT(F9.2)
END

```

```

SUBROUTINE GLOBE(KN,W,T,CTS,CTG)
DIMENSION T(18),CH(41,40),W(18),CTS(20),CTG(20)
COMMON//E1(41,40),E(41,40)
S,PHI(40),PHIH(40),TCG(20,20),QO(18),QOZS(18)
KNP=KN+1
DO 70 I=1,KN
PHI(I)=PFUNC(I,KN)
70 CONTINUE
PHI(KNP)=1.
DO 71 I=1,KN
PHIH(I)=PHFUNC(I,PHI)
71 CONTINUE
PHIH(KNP)=1.
CALL COLUZ(KN,W,T)
DO 72 L=1,KNP
PRINT 25, PHI(L),PHIH(L),W(L),QO(L),QOZS(L)
72 CONTINUE
CALL CMQZ96(KN,CH,CTS,CTG)
25 FORMAT(5E15.5)
5 FORMAT(E9.2)
15 FORMAT(9F8.3)
RETURN
END

```

```

SUBROUTINE COLOZ(KN,W,T)
COMMON/E1(41,40),E(41,40)
PHI(40),PHIH(40),TCG(20,20),QO(10),QOZS(10)
DIMENSION W(10),T(10)
REAL M3,MAIR
KNP=KN+1
MAIR=28.966
AVG=6.0225E+23
M3=44./AVG
RA=8.3143E+07/MAIR
OSTP=2.6869E+19
ATQOS=1.013E+06
G=981.
C ASSUME OZONE SCALE HEIGHT IS H ABOVE TOP HALF LEVEL
H=4.3E+05
QO(1)=PHIH(1)*ATMOS*W(1)/(M3*OSTP*RA*T(1))
QOZS(1)=QO(1)*H
DO 2 K=2,KNP
  QO(K)=PHIH(K)*ATMOS*W(K)/(M3*OSTP*RA*T(K))
C ASSUME W VARIES LINEARLY WITH PRESSURE
  QOZS(K)=QOZS(K-1)+0.5/(M3*OSTP*G)*(W(K-1)+W(K))*((PHIH(K)-
  *PHIH(K-1))*ATMOS
2 CONTINUE
  RETURN
  END

FUNCTION PFUNC(J,N)
C THIS SUBROUTINE COMPUTES PRESSURE LEVELS
  RJN=J-N
  PFUNC=0.1*EXP(RJN*5./7.)
  IF(J.EQ. 1) PFUNC=0.
  RETURN
  END

FUNCTION PHFUNC(J,PHI)
C THIS SUBROUTINE COMPUTES PRESSURE HALF LEVELS
  DIMENSION PHI(10)
  PHFUNC=SQRT(PHI(J)*PHI(J+1))
  IF(J.EQ. 1) PHFUNC=0.5*PHI(2)
  RETURN
  END

FUNCTION GFUNC(A,B,C)
C PLANCK FUNCTION LINEARLY INTERPOLATED IN LOG PRESSURE.
  GFUNC= ALOG(A/C)/ALOG(B/C)
  RETURN
  END

```

```

SUBROUTINE CHOZ96(N,R,CS,CG)
  DIMENSION T(21),R(41,40),Q(20),CTS(20),THETA(20),CS(20),CG(20)
  1, CTS(20),SG(20)
  COMMON//E1(41,40),E(41,40)
  1, PHI(40),PHIH(40),TCG(20,20),QO(10),QOZS(10)
  REAL NUZ,KB,MASS
  DATA M,M1/1,1/
  DATA T/187.1,186.9,193.2,201.6,210.4,221.0,234.7,248.4,263.5,
  270.7,256.0,243.2,229.6,224.2,219.6,216.7,231.4,288.2,300./
C  WHENRBCIN AMENDMENTS
  PI=3.1415926
  SPI=SQRT(PI)
  HP=6.6257E-27
  N1=N+1
  AVG=6.0225E+23
  MASS=44./AVG
  C=2.8979E+10
  KB=1.3805E-10
  CP=1.E+07
  GRAV=981.
  TZ=300.
C  PARAMETERS FOR 9.6 MICRON BAND
  NUZ=1040.
  D=EXP(HP*C*NUZ/(KB*TZ))-1.
  BBAR=2.*HP*(C*NUZ)**3/(C*D)
  CGHST=2.*PI*BBAR*GRAV/CP
C
C
C  COMPUTE CURTIS-GODSON TEMPERATURE FOR EACH PAIR OF LEVELS
  DO 72 I=1,N1
    DO 73 J=I,N1
      TCG(I,J)=T(I)
      IF(I.EQ. J) GO TO 73
      SUM=0.
      J=1+J-1
      /O 76 L=I,JM1
      SUM=SUM+(T(L)+T(L+1))*(PHIH(L+1)-PHIH(L))
  70 CONTINUE
      TCG(I,J)=0.5*SUM/(PHIH(J)-PHIH(I))
      TCG(J,I)=TCG(I,J)
  73 CONTINUE
      PRINT 100,(TCG(I,J),J=1,N1)
  72 CONTINUE
      PRINT 101, M,M1
      CALL FLUXEL(N,M,M1)
C
C  COMPUTE LTE CURTIS MATRIX
  DO 20 I=1,N
    CS(I)=0.
    UP=1.013E+06*(PHI(I+1)-PHI(I))
    DO 21 J=1,N1
      R(J,I)= E1(J+1,I+1)*E1(J,I)-E1(J+1,I)-E1(J,I+1)
      R(J,I)=(CONST/OP)*R(J,I)*86400.
      CS(I)=CS(I)+R(J,I)
  21 CONTINUE
  20 CONTINUE
C
      PRINT 100,((R(J,I),J=1,N1),I=1,N)
      PRINT 102
      PRINT 100,(T(I),I=1,N1)
      PRINT 102
      PRINT 100, (CS(I),I=1,N)
      PRINT 102
C  COMPUTE COOLING RATES
      VZ=HP*C*NUZ/(KB*TZ)
      DO 16 J=1,N1
        V=HP*C*NUZ/(KB*T(J))
        THETA(J)=(EXP(VZ)-1.)/(EXP(V)-1.)
  16 CONTINUE
      PRINT 105, (THETA(J),J=1,N1)
      PRINT 102
      DO 11 J=1,N
        CTS(J)=CS(J)*THETA(J)
        CG(J)=R(N1,J)
        CTG(J)=CG(J)*(THETA(N1)-THETA(J))

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      SG(J)=CTS(J)+CTG(J)
      CS(I)=(D+1.)*CS(I)
      CG(I)=(D+1.)*CG(I)
11  CONTINUE
      DO 10 I=1,M
      Q(I)=0.
      DO 9 J=1,M1
      Q(I)=Q(I)+R(J,I)*THETA(J)
      9  CONTINUE
      PRINT 100, CTS(I),CTG(I),SG(I),Q(I)
10  CONTINUE
C 99 CONTINUE
C 106 FORMAT(F6.1)
105 FORMAT(E10.3)
103 FORMAT(4(2X,E10.4,2X,E10.4,8X))
104 FORMAT(8F6.0)
110 FORMAT(8F6.3)
C 104 FORMAT(I5)
102 FORMAT(/)
C 107 FORMAT(9E10.3)
100 FORMAT(9F10.3/9F6.3)
108 FORMAT(4E10.3/8E10.3)
101 FORMAT(4I3,F6.1)
      RETURN
      END

      SUBROUTINE FLUXEL(N,M,M1)
C      SUBROUTINE FOR CALCULATING CURTIS COEFFICIENTS FOR NET UPWARD FLUX
      DIMENSION C(41,40)
      COMMON//E1(41,40),E(41,40)
      S,PHI(40),PHIH(40),TCG(20,20),QD(18),QDZS(18)
      N1=N+1
      N2=N+2

C
C
      DO 70 I=1,M1
      DO 50 J=1,M
      IF (J.EQ.1) GO TO 46
      IF (J.EQ.I) GO TO 30
      IF (J.GE.N) GO TO 48

C
      E(J,I)=EFUNC(PHI(J),PHI(I))
      COR=0.

C
      DO 15 K=1,M
      EK1=K-1
      EK=K
      PHIK=((M-EK1)/M)*PHIH(J-1)+(EK1/M)*PHIH(J)
      PHIK1=((M-EK)/M)*PHIH(J-1)+(EK/M)*PHIH(J)
15  COR= COR+(GFUNC(PHIK,PHIH(J),PHIH(J-1))+GFUNC(PHIK1,PHIH(J),
      PHIH(J-1)))+(EFUNC(PHIK1,PHI(I))-EFUNC(PHIK,PHI(I)))

C
      C(J,I)=EFUNC(PHIH(J),PHI(I))-E(J,I)-COR/2.
      GO TO 51

C
30  CORON=0.
      CURUP=0.
      E(I,I)=0.

C
      DO 35 K=1,M1
      EK1=K-1
      EK=K
      PHIK=((M1-EK1)/M1)*PHIH(I-1)+(EK1/M1)*PHI(I)
      PHIK1=((M1-EK)/M1)*PHIH(I-1)+(EK/M1)*PHI(I)
      CGKON= CGKON+(GFUNC(PHIK,PHIH(I),PHIH(I-1))+GFUNC(PHIK1,PHIH(I),
      PHIH(I-1)))+(EFUNC(PHIK,PHI(I))-EFUNC(PHIK1,PHI(I)))
      PHIK=((M1-EK1)/M1)*PHI(I)+(EK1/M1)*PHIH(I)
      PHIK1=((M1-EK)/M1)*PHI(I)+(EK/M1)*PHIH(I)
35  CORUP= CORUP+(GFUNC(PHIK,PHIH(I),PHIH(I-1))+GFUNC(PHIK1,
      PHIH(I),PHIH(I-1)))+(EFUNC(PHIK1,PHI(I))-EFUNC(PHIK,PHI(I)))

C
      C(I,I)=EFUNC(PHIH(I),PHI(I))+(CORON-CORUP)/2.

C
      GO TO 51

```

```

46 C(1,I)=0.
   E(1,I)= EFUNC(0.,PHI(I))
   GO TO 51
C
48 E(N,I)= EFUNC(PHI(N),PHI(I))
   L(N1,I)= EFUNC(PHI(N1),PHI(I))
   COR=0.
   M2= 3*M/2
   DO 49 K=1,M2
   EK1=K-1
   EK=K
   PHIK= ((M2-EK1)/M2)*PHI(N-1)+(EK1/M2)*PHI(N1)
   PHIK1= ((M2-EK)/M2)*PHI(N-1)+(EK/M2)*PHI(N1)
49 COR = COR+(GFUNC(PHIK,PHI(N),PHI(N-1))+GFUNC(PHIK1,PHI(N)
   1,PHI(N-1)))+(EFUNC(PHIK1,PHI(I))-EFUNC(PHIK,PHI(I)))
   C(N,I)= E(N1,I)-E(N,I)-COR/2.
C
51 E1(J,I)=E(J,I)+C(J,I)
50 CONTINUE
C
   E1(N1,I)=EFUNC(PHI(N1),PHI(I))
   E1(N2,I)= EFUNC(-1.,-1.)
70 CONTINUE
C
   PRINT 102
   PRINT 101, ((E(J,I),J=1,N),E1(N1,I),E1(N2,I),I=1,N1)
101 FORMAT('F8.3/10F8.3')
   PRINT 102
100 FORMAT(10E10.3)
102 FORMAT(/)
   RETURN
   END

FUNCTION EFUNC(A,B)
C THIS SUBROUTINE COMPUTES FLUX EQUIVALENT WIDTH
C NOTE THAT NAME OF ABSORPTION FUNCTION APPEARS TWICE
C DIMENSION Y(2)
   COMMON//E1(41,40),E(41,40)
   S,PHI(40),PHI1(40),TCG(20,20),QO(18),QOZS(18)
   DATA Y/0.2113,0.7887/
   EFUNC=0.
   IF(A .GE. 0.) GO TO 4
   EFUNC=QZ96A(-1.,0.,0.)
   RETURN
4 CONTINUE
C IA IS PRESSURE LEVEL CLOSEST TO A, IAU IS LEVEL JUST ABOVE A
   IA=IT(A)
   IAU=IU(A)
C SIMILARLY FOR IB AND IBU
   IB=IT(B)
   IBU=IU(B)
   AU=QOZS(IAU)+(A-PHI(IAU))/(PHI(IAU+1)-PHI(IAU))*(QOZS(IAU+1)-
   S QOZS(IAU))
   BU=QOZS(IBU)+(B-PHI(IBU))/(PHI(IBU+1)-PHI(IBU))*(QOZS(IBU+1)-
   S QOZS(IBU))
   U=ABS(AU-BU)
   IF(U .EQ. 0.) GO TO 80
C CUFF15-GOODSON APPROXIMATION
   T=TCG(IA,IB)
   P=(A+B)/2.
C QUADRATURE OVER ANGLES
C TWO POINT GAUSSIAN
   DO 22 L=1,2
   EFUNC=EFUNC+0.5*Y(L)*QZ96A(U/Y(L),T,P)
22 CONTINUE
C
60 CONTINUE
   RETURN
   END

```

```

      FUNCTION IU(X)
C     THIS SUBROUTINE COMPUTES HALF LEVEL NUMBER JUST ABOVE PRESSURE X
      IU=1
      IF(X .LE. 0.) RETURN
      IF(X .GT. 0.1) GO TO 2
      IU=7./5.*ALOG(10./1.4292*X)+17.
      IF(IU .LT. 1) IU=1
      RETURN
2     IU=16
      IF(X .GT. 0.31623) IU=17
      RETURN
      END

```

```

      FUNCTION IT(X)
C     THIS SUBROUTINE COMPUTES HALF LEVEL NUMBER NEAREST PRESSURE X
C     FOR CURTIS-GODSON TEMPERATURE APPROXIMATION
      IT=1
      IF(X .LE. 0.) RETURN
      IF(X .GT. 0.1) GO TO 2
      R=7./5.*ALOG(10./1.4292*X)+17.
      IT=R
      E=R-IT
      IF(E .GT. 0.5) IT=IT+1
      IF(IT .LT. 1) IT=1
2     IU=17
      IF(X .GT. 0.5623) IT=18
      RETURN
      END

```

```

      FUNCTION OZ96A(U,T,P)
C     THIS SUBROUTINE COMPUTES INTEGRATED ABSORPTION OF 9.6 MICRON
C     OZONE BAND USING FORMULA OF AIDA (1975) JQSRT 15, PP. 389-403.
C
      IF(U .GE. 0.) GO TO 6
      OZ96A=260./2.
      RETURN
6     CONTINUE
      PI=3.141592654
      SB=393.4
C     CM-1 CHATH-1 AT 293.2 K
      A=8=0.
      DB=260.
      X=SB*U/DB
      G=0.076*P*(T/293.)**(-0.7)
      D=0.1
      BT=2.*PI*G/D
      Y=2.*X/BT
      E=X/SQRT(1.+Y)
      OZ96A=DB*(1.-EXP(-E))
C
      Z=ALOG10(X/2.82)
C     MODIFICATION TO FORMULA BY WMW
      Z=ABS(Z)
      F=ALOG10(P)
      IF(U .GT. 2.) GO TO 1
      A=0.16524*F**4+0.48623*F**3+0.54604*F**2-0.11218*F+0.04140
      A=10**A
      B=(-0.00412*F**3-0.08904*F**2+0.14168*F)*Z-0.04420*F**3-0.04262*
      F**2+0.50158*F+2.1600
      GO TO 2
1     A=-0.05097*F**3+0.11180*F**2-0.21613*F+0.21748
      A=10**A
      B=-0.4800*F*Z+1.720
2     C=-0.01285*F**4-0.07222*F**3-0.10718*F**2+0.04269*F+0.55000
C
      CF=1.-C*EXP(-(Z**B)/A)
      OZ96A=OZ96A*CF
      RETURN
C
C     ENTRY OZ96W
C
C     THIS SUBROUTINE COMPUTES INTEGRATED ABSORPTION OF 9.6 MICRON
C     OZCNE BAND USING FORMULA OF WALSHAW (1957) QJRM 83, PP. 315-321
C
      IF(U .GE. 0.) GO TO 7
      OZ96A=138./2.
      RETURN
7     CONTINUE
      X=(1.+0.1025*U)/(1.+1.61*U)
      IF(U .LE. 0.1) GO TO 4
      X=0.984*10**(-0.53*U)
      IF(U .LE. 0.4) GO TO 4
      X=0.317*U**(-0.74)
4     A=2.11
      G=U*(X**A)/(760.*P)
      GL=ALOG10(G)
C
C     TABULATED FUNCTION FITTED WITH STRAIGHT LINE SEGMENTS
      E=1.
      IF(GL .LT. -4.2) GO TO 3
      E=-0.0176*(GL+3.8)+0.992
      IF(GL .LT. -2.8) GO TO 3
      E=0.1348*(GL+3.0)+0.948
3     F=1.165/SQRT(1.+734.*G)*E
      OZ96A=138.*(1.-10**(-U*X*F))
      RETURN
      END

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	SUBROUTINE OXY(T,KN)	OXY
C	SUBROUTINE TO COMPUTE DENSITY AND COLUMN ABUNDANCE OF	OXY
C	MOLECULAR OXYGEN	OXY
C		OXY
C	ASSUME N(O2)=0.2095*N(TOTAL) TO 155 KM, AND N(O2)=1.E+17*EXP(-2	OXY
C	5/30 KM) ABOVE 155 KM	OXY
	COMMON/O2/OO2(17),CA2(17)	OXY
	DIMENSION T(17)	OXY
	REAL KN	OXY
	KNL=KN-1	OXY
	KB=1.3805E-16	OXY
	DO 2 J=1,KN	OXY
	RJN=1-J	OXY
	P=0.06996*EXP(RJN*5./7.)	OXY
	R1N=1-KNL	OXY
	IF(J.EQ. KNL) P=0.5*0.1*EXP(R1N*5./7.)	OXY
	IF(J.EQ. KN) P=0.1	OXY
	OO2(J)=0.2095*1.013E+06*P/(KB*T(J))	OXY
	CA2(J)=1.711E+15+4.496E+24*(P-3.645E-09)	OXY
2	CONTINUE	OXY
	RETURN	OXY
	END	OXY
	FUNCTION O2UV(U)	O2UV
C	SUBROUTINE TO CALCULATE SOLAR ENERGY ABSORBED BY MOLECULAR OXYGEN	O2UV
C	PARAMETERIZATION BY STRUBELL (1973) J GEOPHYS RES 63 PP. 6223-6230	O2UV
C		O2UV
C	SCHUMANN-RUNGE CONTINUUM, 1250-1520 A	O2UV
	EX1=0.	O2UV
	E1=-1.E-17*U	O2UV
	IF(E1.LT. -100.) GO TO 3	O2UV
	EX1=EXP(E1)	O2UV
3	EX1=1.1E-17*EX1	O2UV
C	SCHUMANN-RUNGE CONTINUUM, 1520-1750 A	O2UV
	EX2=0.	O2UV
	E2=-2.9E-19*U	O2UV
	IF(E2.LT. -100.) GO TO 4	O2UV
	EX2=EXP(E2)	O2UV
4	EX3=0.	O2UV
	E3=-1.54E-18*U	O2UV
	IF(E3.LT. -100.) GO TO 6	O2UV
	EX3=EXP(E3)	O2UV
6	EX4=0.	O2UV
	E4=-1.1E-17*U	O2UV
	IF(E4.LT. -100.) GO TO 7	O2UV
	EX4=EXP(E4)	O2UV
7	Q2=(3.43*(EX2-2.08*EX3-1.35*EX4)/U	O2UV
C	SCHUMANN-RUNGE BANDS	O2UV
	Q3=9.03E-19	O2UV
	IF(U.LT. 1.E+13) GO TO 2	O2UV
	Q3=1./(0.143*U+9.64E+00*SQRT(U))	O2UV
2	O2UV=Q1+Q2+Q3	O2UV
	RETURN	O2UV
	END	O2UV
	FUNCTION O3UV(U)	O3UV
C	SUBROUTINE TO CALCULATE SOLAR ENERGY ABSORBED BY OZONE	O3UV
C	(DERIVATIVE OF LACIS-HANSEN EXPRESSION)	O3UV
C		O3UV
	F1 = 1.0+(0.042*U)+(3.23E-4*(U**2.0))	O3UV
	F2 = (0.0212/F1)*(1.0-((U/F1)*(0.042+(6.46E-4*U))))	O3UV
	F1 = 1.0+(138.6*U)	O3UV
	F2 = F2+(1.082/(F1**0.805))*(1.0-((138.6*0.805*U)/F1))	O3UV
	F1 = 1.0+((103.6*U)**3.0)	O3UV
	O3UV = F2+((0.0658/F1)*(1.0-(3.0*(F1-1.0)/F1)))	O3UV
	RETURN	O3UV
	END	O3UV

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SUBROUTINE HEAT(TIME,GBV,PB,TB,QB,KHZ,DENS,DZ,JM,KN,TZO,BVF,SH,PI,
$STRODAY,DELTIM,PHO)
COMMON/BUMLG/TG(18),TROP(18),PRS(7,19),QOG
$(18),QU(18,17),COZSG(18),QNZS(18,17),CS(18,17),CG(18,17),ZN(18),
$AZN(18),FUF(18),CZF(19,17),SZF(19,17),COT(19,17)
COMMON/CURTIS/PJI(16),PHIH(16),G1(17),G2(16),E(6,19,18)
$,R(36,17,18)
DIMENSION GBV(KN), PB(JM,KN), TB(JM,KN), QB(JM,KN), DENS(KN), TZO
$(KN),TP(18),TS(18),CM(2,18,17)
DATA TS/216.7,219.6,224.2,229.6,243.2,258.0,270.7,263.5,248.4,
$234.7,221.0,210.4,201.6,193.2,186.9,187.1,231.4,286.2/
JML = JM-1
KNL = KN-1
KNP=KN+1 $ KN2=KN+2
J2=JML/2
C COMPUTE ZONAL MEAN TEMPERATURE (DEVIATION FROM GLOBAL AVERAGE)
DO 65 J=1,JML
DO 65 K=1,KNL
65 TB(J,K)=KHZ+(PB(J,K+1)*DENS(K+1)-PB(J,K)*DENS(K))/DZ
DO 66 J=1,JML
DO 67 K=1,KNL
TP(K)=TB(J,K)+TZO(K)
67 CONTINUE
C CHOOSE APPROPRIATE TROPOSPHERIC TEMPERATURE
TP(KN)=TROP(J)
C TP(KN)=TZO(KN)
IP(KNP)=TG(J)
C COMPUTE SOLAR GEOMETRY FACTORS
DAY = TIME/(24.*60.*60.)
IDAY=DAY+0.005
RES = DAY-IDAY
EPS=0.7*DELTIM/(24.*60.*60.)
IF(RES.GT.EPS) GO TO 60
C COMPUTATION OF EARTH-SUN DISTANCE (BERGER, 1976: J. ATMOS. SCI.
C 35 PP.2362-2367)
EC=0.016722
BETA=SQRT(1.-EC*EC)
OMB=(101.21972+180.)*PI/180.
RLAMBZ=-2.*((EC/2.+(EC**3)/8.)*(1.+BETA)*SIN(-OMB)
$-EC*EC/4.*(0.5+BETA)*SIN(-2.*OMB)
$+(EC**3)/8.*(1./3.+BETA)*SIN(-3.*OMB))
C RLAMBZ IS LONGITUDINAL POSITION AT 21 MARCH (VERNAL EQUINOX)
C ONE YEAR HAS 360 DAYS
RLAMB=RLAMBZ+DAY/360.*2.*PI
RLAMB=RLAMB+(2.*EC-0.25*EC**3)*SIN(RLAMB-OMB)
$+(5./4.)*EC*EC*SIN(2.*RLAMB-2.*OMB)
$+(13./12.)*EC**3*SIN(3.*RLAMB-3.*OMB)
V=RLAMB-OMB
RHO = (1.-EC*EC)/(1.+EC*COS(V))
C SEASONAL VARIATION OF SOLAR HEATING (COGLEY AND BORUCKI (1976)
C J. ATMOS. SCI. 33 PP. 1347-1356, APPROXIMATION 1)
C DELTA=SOLAR DECLINATION
DELTA = 0.4093198*SIN(2.*PI*(STRODAY+DAY-80.)/360.)
DO 30 L=1,JML
PHD=(J2+1-L)*10-5
C PHI=TERRESTRIAL LATITUDE
PHI = PHD*PI/180.
C TSTAR=TIME OF SUNSET (OR NEGATIVE TIME OF SUNRISE)
C ZEN=AVERAGE VALUE OF COS(SZA) BETWEEN -TSTAR AND TSTAR
SUN = TAN(DELTA)*TAN(PHI)
ISUN = SUN
IF (ISUN) 10,15,20
10 TSTAR = 0.
ZEN = COS(DELTA-PHI)
GO TO 25
C
15 TSTAR = (12./PI)*ACOS(-SUN)
ZEN = SIN(DELTA)*SIN(PHI)+(12./(PI*TSTAR))*COS(DELTA)*COS(PHI)*
$ SIN(PI*TSTAR/12.)
GO TO 25
C
20 TSTAR = 12.
ZEN = SIN(DELTA)*SIN(PHI)
25 CONTINUE

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	SUBROUTINE RADEQU(TZO,RHO,JH,KN)	RADEQU
	DIMENSION TZO(17),A3(2)	RADEQU
	DATA ITOT,DT,EPS/20,0.1,0.1/	RADEQU
	JHL=JH-1 3 KNL=KN-1	RADEQU
	PI=3.141592654	RADEQU
	DO 40 IT=1,ITOT	RADEQU
	SUM = 0.	RADEQU
C	IMAX IS KNL FOR CLIMATOLOGICAL TROPOSPHERIC TEMPERATURE,	RADEQU
C	KN FOR RADIATIVE EQUILIBRIUM TROPOSPHERIC TEMPERATURE	RADEQU
	IMAX=KNL	RADEQU
	DO 35 I=1,IMAX	RADEQU
	TZO(I) = TZO(I)+2.*DT	RADEQU
	DO 30 J=1,2	RADEQU
	TZO(I) = TZO(I)-DT	RADEQU
	A3(J) = 0.	RADEQU
	DO 25 L=1,JHL	RADEQU
	THETA = 175-(L-1)*10	RADEQU
	THETA = THETA*PI/180.	RADEQU
	W = SIN(THETA)*PI/30.	RADEQU
	A3(J)=A3(J)+DELT(TZO,L,I,RHO,KN)*W	RADEQU
25	CONTINUE	RADEQU
30	CONTINUE	RADEQU
	DQ = (A3(1)-A3(2))/DT	RADEQU
	DIFF = A3(2)/DQ	RADEQU
	TN = TZO(I)-DIFF	RADEQU
	DTP = DIFF	RADEQU
	IF (DTP.GT.20.) DTP = 20.	RADEQU
	IF (DTP.LT.-20.) DTP = -20.	RADEQU
	TN = TZO(I)-DTP	RADEQU
	DIFF = ABS(DIFF)	RADEQU
	IF (DIFF.GT.SUM) SUM = DIFF	RADEQU
	TZO(I) = TN	RADEQU
35	CONTINUE	RADEQU
	IF (SUM.LT.EPS) GO TO 45	RADEQU
40	CONTINUE	RADEQU
	PRINT 65, ITOT	RADEQU
	STOP 1	RADEQU
45	PRINT 70, IT	RADEQU
	WRITE(6,75) TZO(KN),(TZO(K),K=1,KNL)	RADEQU
	RETURN	RADEQU
C		RADEQU
65	FORMAT (5X,*TEMPERATURE PROFILE FAILED TO CONVERGE AFTER*,I3,	RADEQU
	5* ITERATIONS*)	RADEQU
70	FORMAT (5X,*TEMPERATURE CONVERGED AFTER *,I2,* ITERATIONS*)	RADEQU
75	FORMAT (1X,18F7.2)	RADEQU
	END	RADEQU

	FUNCTION DELT(TP,L,I,RHO,KN)	DELT
	COMMON/BUHLG/TG(L),TRUP(L),PRS(7,19),QOG	DELT
	S(18),OO(18,17),QJZSG(18),QOZS(18,17),CS(18,17),CG(18,17),ZN(18),	DELT
	SAZN(18),FOY(18),CZT(19,17),SZT(19,17),COT(19,17)	DELT
	COMMON/CURTIS/PHI(18),PHIH(18),G1(17),G2(18),E(6,19,18)	DELT
	S,R(30,17,18)	DELT
	COMMON/Q2/UQ2(17),CA2(17)	DELT
	DIMENSION TP(17),TD(18)	DELT
	KNL=KN-1 & KNP=KN+1	DELT
	JML=18	DELT
C	WHW AMENDMENTS, 20 OCTOBER, 1980	DELT
C	V=H+C*WN/K8	DELT
	V15=959.96 & V10=1409.64	DELT
	C15=C10=0.	DELT
C	CHOOSE APPROPRIATE TROPOSPHERIC TEMPERATURE	DELT
	TD(1)=TG(L)	DELT
	TD(2)=TRUP(L)	DELT
C	TD(2)=TP(KN)	DELT
	DO 4 J=3,KNP	DELT
	TD(J)=TP(J-2)	DELT
4	CONTINUE	DELT
	DO 3 J=1,KNP	DELT
	TT=TD(J)	DELT
	TH15=(EXP(V15/300.)-1.)/(EXP(V15/TT)-1.)	DELT
	TH10=(EXP(V10/300.)-1.)/(EXP(V10/TT)-1.)	DELT
	C15=C15+R(L+JML,I,J)*TH15	DELT
	C10=C10+R(L, I,J)*TH10	DELT
3	CONTINUE	DELT
	CCO2=C15+C10	DELT
	COT(L,I)=CCO2	DELT
	CTS=CS(L,I)*EXP(-1500./TP(I))	DELT
	CTG=CG(L,I)*(EXP(-1500./TG(L))-EXP(-1500./TP(I)))	DELT
	COZONE=CTS+CTG	DELT
	CZT(L,I) = COZONE	DELT
	DELT=CCO2+COZONE	DELT
	SO=0.	DELT
	SZT(L,I)=SO	DELT
	ZEN=ZN(L)	DELT
	AZEN=AZN(L)	DELT
	FDAY=FDY(L)	DELT
	IF(FDAY .LT. 0.0001) GO TO 25	DELT
	DZ = 5000.	DELT
	SH = 7000.	DELT
	Z = (1-I)*DZ	DELT
	AIRDEN=1.013E+06*0.06996*EXP(-Z/SH)/(2.87E+06*TP(I))	DELT
	IF(I .EQ. KN) AIRDEN=1.103E+06*0.3162/(2.87E+06*TD(2))	DELT
	TAU = 10.	DELT
	CLOUD = 0.446	DELT
	RG = 0.5	DELT
	ICLOUD=KN	DELT
C	HEATING BY THE ABSORPTION OF SOLAR RADIATION BY OZONE	DELT
C	FROM LACIS AND HANSEN, J ATMOS SCI 31 118-133	DELT
C	CLEAR SKY	DELT
	RAB = 0.219/(1.+0.816*ZEN)	DELT
	RDB = 0.144	DELT
	RBI = RAB+(1.-RAB)*(1.-RDB)*RG/(1.-RDB*RG)	DELT
	KB = RBI	DELT
C	CLOUDY SKY	DELT
	RAB = 0.13*TAU/(1.+0.13*TAU)	DELT
	RDB = RAB	DELT
	RBI = RAB+(1.-RAB)*(1.-RDB)*RG/(1.-RDB*RG)	DELT
	UT = AZEN*QOZS(L,ICLOUD)+1.9*(QOZS(L,ICLOUD)-QOZS(L,I))	DELT
	A1 = QZUV(UT)	DELT
	UT = AZEN*QOZSG(L)+1.9*(QOZSG(L)-QOZS(L,I))	DELT
	A2 = QZUV(UT)	DELT
	U = AZEN*QOZS(L,I)	DELT
	A3 = QZUV(U)	DELT
	CP=1.E+07 & SOLFLX=0.1365E+07	DELT
	SG3=FDAY*86400.*SOLFLX*ZEN*QO(L,I)/(CP*AIRDEN)*(AZEN*A3+	DELT
	SCLOUD*1.9*RBI*A1+(1.-CLOUD)*1.9*RBI*A2)	DELT

U=ALEN*CA2(1)
A=D2UV(U)
SU2=A*DO2(1)*85400.*FOAY/(AIRDEN*CP)
SO=(SO3+SU2)/(RHJ+RHO)
SZTIL,I)=SO
25 DELT=DcLT+SO
RETURN
END

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(Changes in main program)

READ(2,193)((F(M,J,I),J=1,KN2),I=1,KNP),M=1,6)	WAVE2
193 FORMAT(13E20.13/),4E20.13)	WAVE2
READ(2,194)(KAP(I),I=1,KN)	WAVE2
READ(2,200) TZO(KN), (TZO(I), I=1,KNL)	WAVE2
191 FORMAT(F7.4)	WAVE2
READ(2,190)((PRS(I,J),I=1,4),J=1,19)	WAVE2
READ(2,195) QOG, (QO(L,KN),L=1,JML), ((QO(L,K),L=1,JML),K=1,KNL)	WAVE2
READ(2,195) QOZSG, (QOZS(L,KN),L=1,JML), ((QOZS(L,K),L=1,JML),	WAVE2
SK=1,KNL)	WAVE2
READ(2,195) (CTS(L,KN),L=1,JML), ((CTS(L,K),L=1,JML),K=1,KNL)	WAVE2
READ(2,195) (CG(L,KN),L=1,JML), ((CG(L,K),L=1,JML),K=1,KNL)	WAVE2
180 FOMKAT (4F15.4)	WAVE2
195 FORMAT (9E10.3)	WAVE2
200 FORMAT (9F10.3)	WAVE2
DO 20 K=1,KN	WAVE2
IF (K.GT,KNL) GO TO 20	WAVE2
C NEWTONIAN COOLING COEFFICIENTS FOR EDDIES	WAVE2
KAP(K)=KAP(K)*DT/86400.	WAVE2
20 CONTINUE	WAVE2
IRAO = 0	WAVE2
IRCT=12	FEB3
C COMPUTE ZONAL MEAN TEMPERATURE (DEVIATION FROM GLOBAL AVERAGE)	WAVE2
DO 64 J=1,JML	WAVE2
DO 64 K=1,KNL	WAVE2
64 TB(J,K)=RHZ*(PB(J,K+1)*DENS(K+1)-PB(J,K)*DENS(K))/DZ	WAVE2
DO 66 J=1,JML	WAVE2
DO 67 K=1,KNL	WAVE2
TP(K)=TB(J,K)+TZO(K)	WAVE2
67 CONTINUE	WAVE2
C CHOOSE APPROPRIATE TROPOSPHERIC TEMPERATURE	WAVE2
TP(KN)=TRUP(J)	WAVE2
C TP(KN)=TZO(KN)	WAVE2
TP(KNP)=TG(J)	WAVE2
C COMPUTE TEMPERATURE-CORRECTED CURTIS MATRICES FOR EACH ZONE	WAVE2
CALL CURT(TP,KN,TS,CMAT)	WAVE2
C CONVERT TO JPH VERTICAL COORDINATE LABEL	WAVE2
DO 68 JJ=1,KNP	WAVE2
DO 69 II=1,KNL	WAVE2
CURMAT(J,II,JJ)=CMAT(1,KN2-JJ,KN-II)	WAVE2
CURMAT(J+JML,II,JJ)=CMAT(2,KN2-JJ,KN-II)	WAVE2
69 CONTINUE	WAVE2
CURMAT(J,KN,JJ)=CMAT(1,KN2-JJ,KN)	WAVE2
CURMAT(J+JML,KN,JJ)=CMAT(2,KN2-JJ,KN)	WAVE2
68 CONTINUE	WAVE2
65 CONTINUE	WAVE2
180 FORMAT (4F15.4)	WAVE2
195 FORMAT (9E10.3)	WAVE2
200 FORMAT (9F10.3)	WAVE2

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